Harnessing Quantum Information Science for Enhancing Sensors in Harsh Fossil Energy Environments



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Outline

- Short introduction
- Why use quantum information science for sensors?
- Predictive quantum simulations for candidate materials
- Preliminary work with quantum control calculations
- What's next?

Introduction

- B.S. in Physics and Chemistry *Rice University* (2001)
- Ph.D. in Physical Chemistry *M.I.T.* (2007)
- Staff Scientist Sandia National Labs (2007-2013)
- Associate Prof. at UC Riverside Physics/MatSci/Chemistry/ChemE (2014-now)
- Visit us (virtually) at <u>http://www.bmwong-group.com</u>



UC Riverside (UCR)

- Official Hispanic Serving Institution
- Demographics:
- 57% first-generation students to attend college
- Designated as *"top-performing institution for African American & Latino/a students"* by The Education
 Trust <u>1 of only 3 institutions in the nation</u>



41.5% | Hispanic or Latino
33.8% | Asian
11% | White
5.6% | Two or More Races
3.4% | International
3.3% | Black or African American
1.1% | Unknown
0.2% | Native Hawaiian or Other Pacific Islander
0.1% | Native American or Alaskan Native

General Project Objectives



NV-Center Sensors

- Nitrogen vacancy (NV) centers: structural point defects in bulk carbon
- Contain stable, localized electron spin that can be used as sensor

• Coherence signals can persist at 700 – 1000 K (essential for harsh fossil energy environments)

• Can be controlled with electromagnetic fields



NV-Center Sensors (cont.)

- NV centers near the surface have not been thoroughly explored
 - Defects at surface can enable sensitive detection of chemical analytes in fossil energy infrastructures (discussed later)



DiVincenzo's Criteria

- DiVincenzo outlined 3 necessary conditions for quantum sensor
 - (1) Must have discrete resolvable energy levels separated by finite transition energy
 - (2) Must be possible to initialize sensor into well-known state and read out
 - (3) Can be coherently manipulated, leading to transitions between energy levels

DiVincenzo's Criteria (cont.)

- Approaches in this project obey all 3 DiVincenzo's criteria:
 - (1) Electron spin in NV center can be excited to quantized energy states \checkmark
 - (2) Electrons in NV centers can be initialized with electromagnetic pulse, which can be simulated with quantum control algorithms ✓
 - (3) NV center spin state has long coherence time with added advantage of sub-nanometer spatial resolution ✓

Near-Surface NV-Centers

• Current resolution of NV-center sensors ~(5 nm)³ (size of large protein)

• Dipolar magnetic field
$$B_{dip} = \frac{\mu\mu_0}{4\pi} \frac{\sqrt{3\cos^2\theta + 1}}{r^3}$$

• Since $B_{dip} \sim \frac{1}{r^3}$, sensitivity can be increased 3 orders of magnitude by reducing distance of NV center from surface by factor of 10



Initial NV-Center Configurations

• Use DFT to down-select initial NV-center configurations



examples of NV-center configurations near top surface of lattice

• Carry out ab initio MD at various temperatures to test their stability

Ground vs. Excited-State QM

• Ground-state QM can do this:



Z. Ali & B. Wong, Nature Comm. 9, 4733 (2018)



J. Guo & B. Wong, ACS Nano 12, 9775 (2018)



B. Wong & J. Azoulay, Science Advances 5, eaav2336 (2019)

• Ground-state QM *cannot do this*:



B. Wong & G. Scholes, PNAS 117, 11289 (2020)

Real-Time Electron Evolution



N. Ilawe & B. Wong, J. Mat. Chem. C 6, 5857 (2018)



C. Lian & B. Wong, J. Phys. Chem. Lett. 7, 4340 (2019)

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Excited-State QM for Dynamics

- (1) NV-center configurations down-selected with DFT
- (2) Excited-state QM will probe *real-time* interactions between NV centers & EM fields to understand sensor mechanisms
- Electromagnetic radiation (i.e., light) has two components
 - Magnetic field (**B**)
 - Electric field (E)



Related Work on Excited State QM

- Can we *control polarization* switching with optical pulses?
- Excellent application of our periodic RT-TDDFT approach
- Focus on BaTiO₃ as prototypical example
- Applied laser pulse:

$$\mathbf{E}(t) = \mathbf{E}_0 \cos(\omega t) \exp\left[-\frac{(t-t_0)^2}{2\sigma^2}\right]$$



Polarization Switching

 With *ρ*(*t*) converged, ionic movement approximated via Hellmann-Feynman forces using Ehrenfest Dynamics

fluence (MW/cm²) 21.94 76.78 87.75

> We achieve polarization switching (*P changes sign*) when fluence = 87.75 MW/cm²

C. Lian & B. Wong, J. Phys.Chem. Lett. 10, 3402 (2019) ¹⁵

Ferroelectric Summary

• Simplified take-home message:

laser-tune PES to change polarization

laser can reversibly switch polarization back!

C. Lian & B. Wong, J. Phys. Chem. Lett. 10, 3402 (2019)

Optimal Control Fields

- Excited-state QM is an initial value problem
- Can we ask the inverse question: "Can we construct fields that *enable desired behavior in NV center*?"

Computer Physics Communications Volume 258, January 2021, 107541

NIC-CAGE: An open-source software package for predicting optimal control fields in photo-excited chemical systems *, **

NIC-CAGE: Novel Implementation of Constrained Calculations for Automated Generation of Excitations

A. Raza. C. Hong, X. Wang, A. Kumar, C. R. Shelton, B. M. Wong, *Comput. Phys. Commun.* **258**, 107541 (2021)

NIC-CAGE Program

- Calculates fields that enables transition to desired final state $|\psi_f\rangle$
- Uses scheme from GRAdient Pulse Engineering (GRAPE) algorithm

temporal shape of E(t)

<u>vertical arrows</u> = gradients indicating how amplitude changes to maximize transition probability

NIC-CAGE Program (cont.)

• Propagate TDSE to get final state $\psi_{N-1}(x, t = T)$ and maximize **J**:

$$\boldsymbol{J}[\psi_{N-1},\boldsymbol{\epsilon}] = \left| \left\langle \psi_f \left| \psi_{N-1} \right\rangle \right|^2 - \int_0^T \boldsymbol{\alpha} \cdot \boldsymbol{\epsilon}(t)^2 dt \quad \boldsymbol{\leftarrow} \quad \text{fluence penalty}$$

• NIC-CAGE provides *new analytic derivatives* of $J[\psi_{N-1}, \epsilon]$:

$$\frac{d\boldsymbol{J}}{d\epsilon_{j+1/2}} \propto \left(\prod_{k=j+1} \mathcal{H}_{+,k}^{-1} \cdot \mathcal{H}_{-,k}\right) \cdot \mathcal{H}_{+,j}^{-1} \cdot \left[\boldsymbol{\mu} \odot \left(\boldsymbol{\psi}_{j+1} + \boldsymbol{\psi}_{j}\right)\right] \quad \text{with} \quad \mathcal{H}_{\pm,n} = \mathbb{I} \pm \mathcal{H}\left[\boldsymbol{x}, \left(n + \frac{1}{2}\right)\tau\right]$$

Take-home message: no matrix exponentials!
 → Much faster than Octopus and QuTiP software packages

NIC-CAGE Examples

• Single anharmonic potential well

red: target $|\psi_1^2|$ blue: NIC-CAGE propagation

A. Raza & B. Wong, *Comput. Phys. Commun.* **258**, 107541 (2021) ²⁰

NIC-CAGE Examples (cont.)

• Single anharmonic potential well

red: target $|\psi_5^2|$ blue: NIC-CAGE propagation

A. Raza & B. Wong, *Comput. Phys. Commun.* **258**, 107541 (2021) ²¹

NIC-CAGE Examples (cont.)

• Anharmonic double-well potential (*restricted propagation time*)

red: target $|\psi_1^2|$ blue: NIC-CAGE propagation

A. Raza & B. Wong, *Comput. Phys. Commun.* **258**, 107541 (2021) ²²

"Closing the Loop"

Conclusion & Acknowledgements

- Predictive quantum simulations provide rational guidance for constructing quantum sensors for fossil energy infrastructures
- Quantum information science *almost perfect application of excited-state quantum calculations*

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Acknowledgment: This material is based upon work supported by the Department of Energy Award Number DE-FE0031896.

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