Pore scale modeling and interpretation of borehole NMR logs



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SUMMARY

Multiphase fluid flow through porous media is of interest to the oil and gas industry and geologic carbon capture and sequestration (CCS) projects. In most of these applications, subsurface flow simulations are used to manage resources and projects by predicting responses due to driving forces such including production and injection activities. In most instances, the only practical way to model complex systems is at the continuum scale where physical (hydrologic) flow properties are aggregated over large volumes. This imposes several complications affecting model-based management. Geologic media is inherently inhomogeneous and anisotropic-properties which can be difficult to simulate accurately at large scales. Confounding this, robust measurements of relative permeability are expensive and time consuming to perform; this generally limits sampling to a few point measurements, which may or may not be characteristic. Geophysical characterization techniques are promising means by which to improve our understanding of multiphase systems. In particular nuclear magnetic resonance (NMR) methods provide signal that is directly proportional to the amount of hydrogen in pore fluids. Additionally, since NMR methods are sensitive to the surface area to volume ratio of pores, it is often possible to estimate hydraulic flow properties from the data. Borehole NMR logs provide in situ estimates of porosity and permeability, however these logs need to be calibrated in order to ensure the accuracy of the estimates. This is particularity important if more than one fluid phase are present. In this poster we compare borehole NMR measurements of multiphase systems with random walk simulations of NMR responses at the pore scale. The simulations provide a powerful interpretation tool for the data and provide insight into the fluid distribution. Additionally, anisotropy and inhomogeneity can easily be investigated in computer simulations.

NMR RANDOM WALK SIMULATION



The simplified NMR equation of motion (1), as well as the SDR relation, are appropriate only on the macroscopic level and are not applicable to pore scale NMR. Instead, stochastic random walk simulations where individual spins within a pore matrix are tracked as they diffuse due to Brownian motion can be used to simulate the NMR response of media at this scale. Single phase NMR code developed by the Imperial College London (ICL) [2] was extended for multiple phases and used to perform synthetic NMR simulations on the extracted pore matrices (a), the simulated data are shown in (b) for a single subcube of $100 \times 100 \times 100$ voxels. The simulations shown in this block are single phase; the permafrost example demonstrates the multiple phase NMR simulation capabilities.

Random walk simulations of subsets of the entire image volume were similarily generated. The results were then stiched together to form a composite simulated NMR image of the core with the pore space filled 100% with water. One slice of the synthetic NMR image is shown including estimates of NMR porosity, log mean T_2 and uncalibrated permeability.

PORE MATRIX FROM μ **CT IMAGES**









0.45

0.40

0.35

> 0.30

P 0.20

freq' freq' Threshold



(b)

A μ CT image (at 11.5 μ m) voxel resolution) of the pore Morrow B Sandstone unit is shown (a). In order to extract a pore matrix suitable for simulation (b) appropriate threshold values must be grair determined. Since image intensity varies throughout the volume, a global threshold value is not appropriate. Instead, subsets of the image are analysed as shown in (c,d). The 320 \times 320 \times 10 voxel image subset in



CMR LOGS







(1)

NMR DYNAMICS

 0.40_{1}

0.35

0.30

Hydrogen atoms in liquid water develop bulk magnetization moments $\mathbf{M}_{N}^{(0)}$ in a static magnetic field (\mathbf{B}_0) . These moments precess at the Larmor frequency and may be tipped away from their equilibrium position using RF radiation at the same frequency. After the tipping pulse the moments decay back to equilibrium, macroscopically characterized by three parameters T_1 , T_2 , and T_2^* [1].

$$\mathbf{M}_{N}^{(0)}(\mathbf{r},T) = 2n_{H_{2}O}B_{0}\frac{\gamma_{H}^{2}\hbar^{2}}{4k_{B}T}f(\mathbf{r})\hat{\mathbf{B}}_{0}$$

$$\mathbf{M}_{N}(\mathbf{r},t) = \mathbf{M}_{N}^{(0)}(\mathbf{r})\left\{1 - e^{-(t-\tau_{p})/T_{1}(\mathbf{r})} + e^{-(t-\tau_{p})/T_{1}(\mathbf{r})}\cos\left[\theta_{T}(\mathbf{r},\tau_{p})\right]\right\}$$

$$+ e^{-(t-\tau_{p})/T_{2}^{*}(\mathbf{r})} \times \left[\mathbf{M}_{N}^{(0)}(\mathbf{r}) \times \hat{\mathbf{B}}_{T}^{+}(\mathbf{r},t)\right]\sin\left[\theta_{T}(\mathbf{r},\tau_{p})\right]$$

$$\hat{\mathbf{B}}_{0} \qquad \hat{\mathbf{B}}_{0} \qquad \hat{\mathbf{B}}_{0} \qquad \hat{\mathbf{B}}_{0}$$



Where γ_H is the gyromagnetic ratio of hydrogen, f is the porosity of the media, and n_{H_2O} is the number density of water. \mathbf{B}_T^+ is the corotating part of the transmitter field (\mathbf{B}_T) and τ_p is the duration of the tipping pulse which together determine the tip angle θ_T . In magnetic media nuclear spins dephase causing T_2^* to dephase and deviate from T_2 ; refocusing CPMG pulses are used such that $T_2^* \rightarrow T_2$. In porous media the permeability can be related using, for example, the SchlumbergerSchlumberger CMR log at Farnsworth Unit, well 1310-A

ACKNOWLEDGEMENTS

Funding for this project is provided by the U.S. Department of Energy's (DOE) National Energy Technology Laboratory (NETL through the Southwest Regional Partnership on Carbon Sequestration (SWP) under Award No. DE-FC26-05NT42591. Additional support has been provided by site operator Chaparral Energy, L.L.C. and Schlumberger Carbon Services.



REFERENCES







