



# HIGH-FREQUENCY TRANSVERSE COMBUSTION INSTABILITIES IN LOW-NOX GAS TURBINES (DE-FE0031285)

PI – TIMOTHY LIEUWEN COPI – WENTING SUN

VISHAL ACHARYA BENJAMIN EMERSON DAVID WU

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## **MOTIVATION FOR RESEARCH**



- Combustion dynamics are a critical challenge in high efficiency gas turbines
  - Transverse, "screeching" instabilities are increasingly problematic
- <u>Target architecture</u> → Multi-nozzle can combustor configuration with interacting flames

 Focus of the project → High-frequency transverse combustion instabilities in multi-nozzle can combustor configurations.



Realistic multi-nozzle experiments and high frequency modeling capabilities needed



# DISTINGUISHING FEATURES OF TRANSVERSE INSTABILITIES

- Acoustic wave motions perpendicular relative to main flow direction
  - intrinsic non-axisymmetric excitation
  - No direct net mass flow modulation of reactants (although indirect effect can be large, F<sub>TL</sub> path on chart)
- Wave can be spinning or standing; nodal lines can move

O'Connor, L., Acharya, V., Lieuwen, T., "Transverse Combustion Instabilities: Acoustics, Hydrodynamics, and Flame Dynamics", *Progress in Energy and Combustion Sciences*, Vol. 49.







# DISTINGUISHING FEATURES OF HIGH FREQUENCY, TRANSVERSE INSTABILITIES

- Acoustic wavelength of the order of heat release zone extent
  - Flame not compact
    - its spatial distribution matters!
    - e.g., ideal heat release for 1-T mode has spatial integral of zero
- Coupling mechanisms
  - Velocity
  - Equivalence ratio
  - Direct kinetic (pressure) coupling?
    - Negligible at "low frequencies"









## **RESEARCH FOCUS**



<u>Key research questions</u>:

(1) How do the conventional coupling mechanisms from low-frequency translate to high-frequency?

(2) How do coherent structures interact with high-frequency acoustic forcing?

(3) What are the new mechanisms that are of importance at high frequencies and what are their relative roles when compared to the conventional mechanisms?

(4) How does the direct effect of pressure fluctuations influence the thermoacoustic stability of the system?



# **TASKS & CONNECTIONS**

- Task 1: PMP
- Task 2: Experiments on selfexcited transverse instabilities
  - 2.1 Design of Experiment
  - 2.2 Pressure characterization
  - 2.3 Flow and flame characterization
- Task 3: Reduced order modeling for thermoacoustic coupling
  - 3.1 Flame response modeling
  - 3.2 Hydrodynamic stability modeling
  - 3.3 Kinetic coupling mechanism modeling





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## TASK 2.1 PROGRESS (EXPERIMENTAL COMBUSTOR)

D=29 cm combustor

 → capture accurate high frequency acoustics

**RESEARCH TASK 2** 

- Multiple nozzles

   → capture flame-flame
   interactions.
- Optical accessibility using quartz
   → spatio-temporal flow and
   flame characterization
- Ability to vary heat release distribution in radial and azimuthal direction
   → 3 fuel circuits.



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- Acquired several hundred operating condition data points •
  - $\rightarrow$  Preheat T, global equivalence ratio, pilot ratio (PR =



Today: Focus attention on 1-T mode at 1600 Hz

 $10^{-1}$ 

 $10^{-2}$ 

 $10^{-4}$ 

0

0.5

 $\frac{|d|}{d}$  10<sup>-3</sup>



<sup>ṁ</sup>pilot

m<sub>pilot</sub>+m<sub>outer</sub>



-), air mass flow rate.



#### CH\* phase averaged images

- Out-of-phase oscillations between the top and the bottom
- Observed coherent structures convecting downstream





• Pressure decomposition

$$\hat{p}(\theta,t) = \underbrace{F(t)e^{i(\theta+\psi_F-\omega_F t)}}_{CCW} + \underbrace{G(t)e^{-i(\theta-\psi_G+\omega_G t)}}_{CW}$$

• Given multiple signals,

$$\begin{bmatrix} \hat{p}(\theta_{1},t) \\ \hat{p}(\theta_{2},t) \\ \vdots \\ \hat{p}(\theta_{n},t) \end{bmatrix} = P(measure) = M(known) = M(known)$$
 
$$= \begin{bmatrix} e^{i\theta_{1}} & e^{-i\theta_{1}} \\ e^{i\theta_{2}} & e^{-i\theta_{2}} \\ \vdots & \vdots \\ e^{i\theta_{n}} & e^{i\theta_{n}} \end{bmatrix} \underbrace{ \begin{bmatrix} Fe^{i(\psi_{F}-\omega_{F}t)} \\ Ge^{i(\psi_{G}-\omega_{G}t)} \end{bmatrix} }_{=X(unknown)} \rightarrow P = M \cdot X$$

• Least squares fit

$$X = (M^T M)^{-1} M^T P = \begin{bmatrix} F e^{i(\psi_F - \omega_F t)} \\ G e^{i(\psi_G - \omega_G t)} \end{bmatrix}$$



#### **RESEARCH TASK 2** TASK 2.2 EXTRACTED PARAMETERS



• Magnitude of each wave (CW and CCW) and spin ratio (SR):

$$X| = \begin{bmatrix} |F| \\ |G| \end{bmatrix} \rightarrow SR = \frac{|F| - |G|}{|F| + |G|} \quad where \begin{cases} SR = 0: standing mode \\ SR = \pm 1: CCW/CW spinning mode \\ o.w: mixed mode \end{cases}$$

- Phase difference between CW and CCW waves,  $\phi$ :
- Anti-nodal line,  $\theta_a$  (location where pressure mag. is maximum):
- Anti-nodal line velocity,  $\Omega_a = \frac{d\theta_a}{dt}$
- Total pressure magnitude:  $(2F^2 + 2G^2)^{1/2}$

#### RESEARCH TASK 2 TASK 2.2 EXAMPLE- EXTRACTED PARAMETERS



#### Identification of acoustic mode shape- Azimuthal dependency



RESEARCH TASK 2 TASK 2.2 MODAL DYNAMICS



Case 1 (standing dominant)

Case 2 (CW dominant)



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#### RESEARCH TASK 2 TASK 2.2 MODAL DYNAMICS



- For case 1, SR oscillates around zero (standing). Anti-nodal line stays at constant value.
- For case 2, SR mainly stays in -0.6 (CW), but intermittently switches to 0 (standing) or 0.6 (CCW). Anti-nodal line dominantly rotates in CW, but sometimes in CCW.





#### Phase Portrait in Spin Ratio & Phase Difference Space

- Three different operating conditions were investigated.
- Acquire data for 30 seconds (~ 50,000 cycles) at stabilized operating conditions
- 1. Remove data points outside two standard deviation of the mean of  $(2F^2 + 2G^2)^{1/2}$

2. Divide phase space into  $M \times M$  grid. Find all data points satisfying  $|SR - SR_i| < \epsilon_{SR}$  and  $|\phi - \phi_i| < \text{simultaneously}$  (i.e.,  $\epsilon_{SR} = 2/M$  and  $\epsilon_{\phi} = 2\pi/M$ ), denoted as  $(SR_i(t_k), \phi_i(t_k))$ 

3. Identify phase space evolution of points originating from  $(SR_i(t_k), \phi_i(t_k))$  after  $\Delta t$ , given by  $(SR_i(t_k + \Delta t), \phi_i(t_k + \Delta t))$ 

4. Compute ensemble average of  $(SR_i(t_k + \Delta t), \phi_i(t_k + \Delta t)): (\overline{SR}_i(t_k + \Delta t), \overline{\phi}_i(t_k + \Delta t))$ 

5. Plot vector from  $(SR_i(t_k), \phi_i(t_k))$  to  $(\overline{SR}_i(t_k + \Delta t), \overline{\phi}_i(t_k + \Delta t))$ 

6. Repeat these steps for each *SR* and  $\phi$  pair,  $(SR_j(t_k), \phi_j(t_k))$ . Results are plotted for all pairs with >100 realizations.

For the results presented here, we used M = 15 and  $\Delta t = 10$  cycles.



- For case 1, stable attractor appears at the center of the phase portrait.
- CW spiral trajectories.





- For case 1, stable attractor appears at the center of the phase portrait.
- CW spiral trajectories.

















0.01 For case 3, multiple fixed points and ٠ 0.005 saddle points were observed.  $\stackrel{ imes 10^3}{7}$ Most realizations stays at A<sub>ccw</sub> 0.6 0.8 -0.8 -0.6 -0.4 0.2 0.4 3 6  $\mathbf{2}$ 51 4 3 Occurrence rad) 0 -1  $\mathbf{2}$ 1 0 0.01 005 -1 -0.50 0.51 SRCase 3 (CCW dominant)

#### RESEARCH TASK 2 KEY TAKEAWAYS



**Developed multi-nozzle facility with radial/azimuthal fuel staging capabilities** 

Identified range of conditions with 1-T mode, with various dominant standing and spinning wave dominant conditions

Single point pressure measurements show pressure amplitudes with variety of characteristics (relatively constant, varying significantly, etc.)

Developed suite of system identification tools for extraction of nodal line position/velocity, spin ratio, amplitude

Phase portraits show different types of dominant attractors, leading to substantively different pressure sensor measurements

Observed coherent structure from CH\* phase averaged images, suggesting potential mechanism of transverse instability.







# RESEARCH TASK 3 REDUCED ORDER MODELING FOR THERMOACOUSTIC COUPLING

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# RESEARCH TASK 3.1 PROGRESS



- Direct excitation by transverse acoustic velocity disturbances  $F_T$
- Indirect excitation through
  - $F_{TL}$  Excited longitudinal acoustics
  - $F_{T\omega}$  Transverse acoustics directly excites inherent flow instabilities
  - $F_{L\omega}$  Transverse acoustics indirectly excites flow instabilities

#### FOCUS – DIRECT EXCITATION MECHANISM EFFECTS



![](_page_23_Picture_9.jpeg)

## RESEARCH TASK 3.1 PROGRESS RESEARCH QUESTIONS

![](_page_24_Picture_1.jpeg)

- What are the relative stabilities for different highfrequency transverse acoustic modes?
- Stability boundaries?
- How do different nozzles in the combustor contribute to stability/instability for the same acoustic mode? (nozzle position, flame shape)

![](_page_24_Figure_5.jpeg)

Pressure (solid) and velocity (dashed) nodal lines for several lower order transverse modes in a circular duct.

Reproduced from Lieuwen, Unsteady Combustor Physics, 2012

![](_page_24_Picture_8.jpeg)

# RESEARCH TASK 3.1 PROGRESS

![](_page_25_Picture_1.jpeg)

High-frequency  $\rightarrow$  Non-compact flames

- Cannot use global heat release dynamics
- Acoustic energy dynamics

$$\frac{\partial}{\partial t} \int_{V} e_{1} dV = \frac{\gamma - 1}{\gamma p_{0}} \int_{V} (p_{1} \dot{q}_{1}) dV - \int_{S} p_{1} \left( \vec{u}_{1} \cdot d\vec{S} \right)$$

A necessary condition for the Rayleigh Integral:

$$RI = \iint_{t \ V} p_1 \dot{q}_1 dV dt > 0$$

![](_page_25_Picture_8.jpeg)

#### **RESEARCH TASK 3.1 PROGRESS** RAYLEIGH CRITERION MECHANICS FOR COMPACT FLAMES

![](_page_26_Picture_1.jpeg)

- RI  $\rightarrow$  Time average of product of two fluctuating quantities
- Sign of RI depends on relative phasing of fluctuating quantities

![](_page_26_Figure_4.jpeg)

# RESEARCH TASK 3.1 PROGRESS Georgia

![](_page_27_Figure_1.jpeg)

From Lovett, J., and Uznanski, K., Prediction of Combustion Dynamics in a Staged Premixed Combustor, ASME Paper # 2002-GT-30646

- Data shows "instability bands" of a single mode
  - instabilities depend upon convective time

![](_page_27_Picture_5.jpeg)

#### RESEARCH TASK 3.1 PROGRESS PRIOR WORK ON DIRECT EXCITATION (SATTELMAYER & CO-WORKERS)

- Phenomenological argument in 2D, using flame displacement mechanism arguments
- Modes are unconditionally unstable for thermoacoustic coupling
  - No stability bands!

![](_page_28_Figure_4.jpeg)

### **RESEARCH TASK 3.1 PROGRESS** DIRECT EXCITATION MODELING IN THIS DOE PROGRAM

![](_page_29_Picture_1.jpeg)

![](_page_29_Figure_2.jpeg)

Level-set for a thin premixed flame

$$\frac{\partial G}{\partial t} + \vec{u} \cdot \vec{\nabla} G = s_L \left| \vec{\nabla} G \right|$$

$$\succ \text{ Explicit flame dynamics}$$

$$G(r, \theta, z, t) = z - \xi(r, \theta, t)$$

$$\frac{\partial \xi}{\partial t} + u_r \frac{\partial \xi}{\partial r} + \frac{u_\theta}{r} \frac{\partial \xi}{\partial \theta} + s_L \left[ \left( \frac{\partial \xi}{\partial r} \right)^2 + \frac{1}{r^2} \left( \frac{\partial \xi}{\partial \theta} \right)^2 + 1 \right]^{\frac{1}{2}} = u_z$$

- Single-valued flame position
- Flame is a thin interface
- Constant flame speed
- Imposed velocity disturbances
- Normalized Rayleigh Index

$$\Re = \frac{RI}{2E}$$

#### **RESEARCH TASK 3.1 PROGRESS** Georgia Tech **RESULTS – COMBUSTOR CENTERED AXISYMMETRIC FLAME** m = 01.2 m = 0-l = 2l = 3Acharya & Lieuwen, -l = 4-l = 50.8 AIAA SciTech 2018 ₩ 0.6 Radial modes – always 0.4 unstable! 0.2 Mixed modes 0.2 0.9 0.30.40.50.6 0.70.8 1 $\beta_F$ (1,1) mode has islands • (l,m) = (1,1)(l,m) = (1,2)of stability Other modes are m = 2m = 1Sta 0.8 0.8 unstable but with lower 0.6 0.F RI than unstable regions 0. 0.2 of (1,1) mode 0.6 0.7 0.8 0.9 0.3 0.4 0.5 0.3 0.4 0.5 0.6 0.7 0.8 0.9 (l,m) = (2,1)(l,m) = (2,2)1.8 1.4 1.2 0.8 0.6 0.6 0.4 0.4 0.2 0.2

0.2

0.3 0.4 0.5

0.8 0.9

0.6 0.7

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0.3 0.4 0.5 0.6 0.7

#### **RESEARCH TASK 3.1 PROGRESS** RESULTS – STABILITY OF FLAMES OFFSET FROM COMBUSTOR AXIS

- Acharya & Lieuwen, ESSCI 2018
- Flame location relative on mode shape affects RI
  - Nozzle offset from axis implies, local azimuthal acoustic velocity affects RI

#### **Outward radial offset increases stability**

![](_page_31_Figure_5.jpeg)

#### Azimuthal offset increases stability

 $(r_c, \theta_c)$  $(r, \theta)$ 

R

С

 $R_{CB}$ 

![](_page_31_Figure_7.jpeg)

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# RESEARCH TASK 3.1 PROGRESS

- Key Takeaways
  - Axisymmetric Mean Flame
    - Combustor centered flame local acoustic azimuthal velocity does not directly influence RI
    - Offset flame local acoustic azimuthal velocity affects RI indirectly.
  - Radial mode unconditionally unstable
  - Mixed mode (1,1) has islands of stability, other modes relatively unstable
  - Radial and azimuthal offset increases stability

### ON-GOING WORK

What happens when mean flame is non-axisymmetric?

What are the stability sensitivities? How do small changes in different control parameters affect RI?

![](_page_32_Picture_11.jpeg)

# **RESEARCH TASK 3.2 PROGRESS**

#### HYDRODYNAMIC STABILITY MODELING

- Swirling flows and flames in annular geometries exhibit range of topologies.
- Complex flow dynamics strongly affect flame stabilization and shape:
  - Flame aerodynamically stabilized in front of vortex breakdown stagnation feature.
  - Flame anchoring position highly unsteady, in contrast to stabilization at edges/corners.
  - Response of flame to imposed disturbance is materially different.

![](_page_33_Picture_7.jpeg)

(d)

![](_page_33_Picture_8.jpeg)

(b)

(a)

![](_page_33_Picture_9.jpeg)

# **RESEARCH TASK 3.2 PROGRESS**

#### HYDRODYNAMIC STABILITY MODELING

- Swirling flows and flames in annular geometries exhibit range of topologies.
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  - Flame aerodynamically stabilized in front of vortex breakdown stagnation feature.
  - Flame anchoring position highly unsteady, in contrast to stabilization at edges/corners.
  - Response of flame to imposed disturbance is materially different.
- Under what circumstances can such flames exist?
  - Not always observed; flames may blowoff directly without reverting to a "free floating" configuration
  - Flow must have interior stagnation point

![](_page_34_Picture_10.jpeg)

![](_page_34_Figure_11.jpeg)

# RESEARCH TASK 3.2 PROGRESS

![](_page_35_Picture_1.jpeg)

![](_page_35_Figure_2.jpeg)

## RESEARCH TASK 3.2 PROGRESS HYDRODYNAMIC STABILITY MODELING

![](_page_36_Picture_1.jpeg)

Nature of centerbody wake/vortex breakdown changes with geometry, swirl number, and Reynolds number

![](_page_36_Figure_3.jpeg)

## RESEARCH TASK 3.2 PROGRESS TAKEAWAYS & FUTURE WORK

![](_page_37_Picture_1.jpeg)

- Key Takeaways
  - Flame shape and stabilization cannot be predicted reliably without excellent understanding of underlying fluid mechanics.
  - Early results indicate that even simple steady laminar flows can exhibit strong hysteresis and three-dimensionality which drives further complexity at higher *Re*.
  - Much more to come here...

### ON-GOING WORK

- How do flow geometry and confinement influence steady/timeaverage dynamical behavior of nonreacting swirling flow systems?
- How does the dynamical picture change in the presence of a reacting flow?

![](_page_37_Picture_9.jpeg)

## **CONCLUDING REMARKS**

![](_page_38_Picture_1.jpeg)

- Developed a multi-nozzle can combustor for studying high-frequency instabilities
  - Wide range of operating conditions and pressures.
  - Multiple pressure sensors to simultaneously capture time-varying pressure
  - Development of modal identification methods for determining the transverse modes using pressure data
  - Identification of mode behaviors by comparing clockwise and counter-clockwise modes in phase space.
- High-frequency flame response modeling
  - Developed a model for direct excitation of flames to high-frequency acoustic modes: Model helps understand the effect of flame position in can as well as flame shape, on mode stability
  - Additional models for coupling mechanisms unique to high-frequency instabilities are in development
    - Pressure coupling effects
    - High-frequency hydrodynamics and its coupling with acoustics and heat release