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IOWA STATE UNIVERSITY

Kinetic Theory Modeling of Turbulent Multiphase Flow

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A Solution Algorithm for Fluid-Particle Flows Across All Flow Regimes

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- 4. Solution Algorithm
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Introduction

- 1. Fluid-particle flow are common in many energy applications, such as fluidized bed and risers.
- 2. Two-Fluid model is the most widely used in simulating this type of flows. However, the hydrodynamic description is inaccurate when particles are dilute.
- 3. Quadrature-based moment methods (QBMM) can be used to find approximate numerical solutions to the particle kinetic equation, thus model particle motions more accurately. But its explicit nature makes it inefficient when particles are close-packed.
- 4. Our objective is to develop a solution algorithm that combines the best features of the hydrodynamic and QBMM solvers, which can accurately simulate fluid-particle flows across all flow regimes.



Fluid-Particle Flow Governing Equations

Fluid Continuity:

$$\frac{\partial \rho_g \alpha_g}{\partial t} + \boldsymbol{\nabla} \cdot \rho_g \alpha_g \boldsymbol{U}_g = 0$$

Fluid Momentum:

$$egin{aligned} &rac{\partial
ho_g lpha_g oldsymbol{U}_g}{\partial t} +
abla \cdot
ho_g lpha_g oldsymbol{U}_g \otimes oldsymbol{U}_g &=
abla \cdot
ho_g lpha_g oldsymbol{\sigma}_g -
abla p_g +
ho_g lpha_g oldsymbol{\sigma}_g -
ho_p lpha_g oldsymbol{M}_{pg} \ &M_{pg} &= rac{1}{ au_p} (oldsymbol{U}_g - oldsymbol{U}_p) - rac{1}{
ho_p}
abla p_g + rac{
ho_g}{
ho_p}
abla \cdot lpha_g oldsymbol{\sigma}_g \end{aligned}$$

Particle Kinetic:

$$\frac{\partial f\left(\mathbf{v}\right)}{\partial t} + \mathbf{v} \cdot \frac{\partial f\left(\mathbf{v}\right)}{\partial \mathbf{x}} + \frac{\partial}{\partial \mathbf{v}} \cdot f\left(\mathbf{v}\right) \mathbf{A} = \mathbf{S}$$

A represents acceleration due to forces acting on each particle, S represents other possible source terms, e.g. particle collisions.





Fluid-Particle Flow Governing Equations

Particle Moments Transport:

$$\begin{aligned} \frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot \mathbf{F} &= \mathbf{S} \qquad \mathbf{M}_{ijk}^{\gamma} = \int v_1^i v_2^j v_3^j f(\boldsymbol{v}) \, d\boldsymbol{v} \\ \alpha_p &= M_{000}^0, \quad \alpha_p \boldsymbol{U}_p = \begin{bmatrix} M_{100}^1 \\ M_{010}^1 \\ M_{001}^1 \end{bmatrix}, \quad \alpha_p \boldsymbol{U}_p \otimes \boldsymbol{U}_p + \alpha_p \mathbf{P}_p = \begin{bmatrix} M_{200}^2 & M_{110}^2 & M_{101}^2 \\ M_{110}^2 & M_{020}^2 & M_{011}^2 \\ M_{101}^2 & M_{001}^2 \end{bmatrix}. \end{aligned}$$

Particle Continuity:

$$rac{\partial
ho_p lpha_p}{\partial t} +
abla \cdot
ho_p lpha_p oldsymbol{U}_p = 0$$

Particle Momentum:

$$rac{\partial
ho_p lpha_p oldsymbol{U}_p}{\partial t} +
abla \cdot
ho_p lpha_p \left(oldsymbol{U}_p \otimes oldsymbol{U}_p + oldsymbol{P}_p + oldsymbol{G}_p + oldsymbol{Z}_p
ight) =
ho_p lpha_p oldsymbol{g}_p +
ho_p lpha_p oldsymbol{M}_{pg}$$

Particle Particle-Pressure Tensor:

$$\frac{\partial \rho_{\rho} \alpha_{\rho} \mathbf{P}_{\rho}}{\partial t} + \nabla \cdot \rho_{\rho} \alpha_{\rho} \left(U_{\rho} \otimes \mathbf{P}_{\rho} + \mathbf{Q}_{\rho} + \mathbf{H}_{\rho} \right) + \rho_{\rho} \alpha_{\rho} \left[\left(\mathbf{P}_{\rho} + \mathbf{G}_{\rho} \right) \cdot \nabla U_{\rho} + \left(\nabla U_{\rho} \right)^{T} \cdot \left(\mathbf{P}_{\rho} + \mathbf{G}_{\rho} \right) \right] = \rho_{\rho} \alpha_{\rho} \mathbf{E}_{\rho g} + \rho_{\rho} \alpha_{\rho} \mathbf{C}_{\rho}$$

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Particle Kinetic, Collisional and Frictional Flux

Kinetic Flux:
$$U_p \otimes U_p + P_p$$
 $P_p = \Theta_p I - \sigma_p$ $\sigma_p = 2\nu_{p,k}S_p$ $S_p = \frac{1}{2} \left[\nabla U_p + (\nabla U_p)^T - \frac{2}{3} (\nabla \cdot U_p) I \right]$ Collisional Flux: $G_p = [2(1+e)\alpha_p g_0 \Theta_p - \nu_{p,b} \nabla \cdot U_p] I - 2\nu_{p,c} S_p$ Frictional Flux: $Z_p = \frac{Fr}{\rho_p \alpha_p} \frac{(\alpha_p - \alpha_{p,fr,min})^r}{(\alpha_{p,max} - \alpha_p)^s} \left(I - \frac{2\sin\phi}{\|S_p\|} S_p \right)$ Granular Pressure: $p_p = p_{p,k} + p_{p,c} + p_{p,f}$ Particle Viscosity: $\nu_p = \nu_{p,k} + \nu_{p,c} + \nu_{p,f}$ Kinetic and Collisional Heat Flux: $Q_p + H_p = -\frac{2}{3}k_\Theta \nabla \otimes P_p$



Operator-splitting scheme for all flow regimes

10.

$$\frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot h_1 \mathbf{F} + \nabla \cdot (h_2 \mathbf{F} + \mathbf{G} + \mathbf{Z}) = \mathbf{S}$$

$$h_1 = 1 - h_2$$

$$h_2 = \left(\frac{p_{p,c}^* + p_{p,f}}{p_{p,k} + p_{p,c}^* + p_{p,f} + \varepsilon}\right)^p$$

$$p_{p,c}^* = 2(1 + \varepsilon)\rho_p \alpha_p^2 g_0 \Theta_p$$

$$\frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot h_1 \mathbf{F} = \mathbf{0}$$

$$\frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot (h_2 \mathbf{F} + \mathbf{G} + \mathbf{Z}) = \mathbf{S}$$

$$h_1 = \frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot (h_2 \mathbf{F} + \mathbf{G} + \mathbf{Z}) = \mathbf{S}$$

$$h_2 = \frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot h_1 \mathbf{F} = \mathbf{0}$$

$$\frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot (h_2 \mathbf{F} + \mathbf{G} + \mathbf{Z}) = \mathbf{S}$$

$$h_1 = \frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot (h_2 \mathbf{F} + \mathbf{G} + \mathbf{Z}) = \mathbf{S}$$

$$h_2 = \frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot (h_2 \mathbf{F} + \mathbf{G} + \mathbf{Z}) = \mathbf{S}$$

$$h_3 = \frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot (h_2 \mathbf{F} + \mathbf{G} + \mathbf{Z}) = \mathbf{S}$$

Hydrodynamic solver:
$$\frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot (h_2 \mathbf{F} + \mathbf{G} + \mathbf{Z}) = \mathbf{S}$$

Particle Volume Fraction:

$$\frac{\partial \alpha_p}{\partial t} + \boldsymbol{\nabla} \cdot h_2 \alpha_p \boldsymbol{U}_p = 0$$

Particle Velocity :

$$\frac{\partial \alpha_p U_p}{\partial t} + \nabla \cdot \left(h_2 \alpha_p U_p \otimes U_p + \frac{p_p^*}{\rho_p} I - 2\alpha_p \nu_p^* S_p \right) = \alpha_p g + \frac{\alpha_p}{\tau_p} \left(U_g - U_p \right) - \frac{\alpha_p}{\rho_p} \nabla p_g + \alpha_p \rho_g \nabla \cdot \alpha_g \sigma_g.$$
$$p_p^* = h_2 p_{p,k} + p_{p,c} + p_{p,f} \qquad \nu_p^* = h_2 \nu_{p,k} + \nu_{p,c} + \nu_{p,f}$$

Particle Granular Temperature :

$$\frac{3}{2} \left(\frac{\partial \alpha_p \Theta_p}{\partial t} + \nabla \cdot h_2 \alpha_p \Theta_p U_p \right) = \nabla \cdot (\alpha_p k_{\Theta}^{\dagger} \nabla \Theta_p) - \left(\frac{p_p^{\dagger}}{\rho_p} \mathbf{I} - 2\alpha_p \nu_p^{\dagger} \mathbf{S}_p \right) : \nabla U_p - 3 \left(\frac{1 - e^2}{2\tau_c} + \frac{1}{\tau_p} \right) \alpha_p \Theta_p$$
$$p_p^{\dagger} = h_2 p_{p,k} + p_{p,c} \qquad \nu_p^{\dagger} = h_2 \nu_{p,k} + \nu_{p,c} \qquad k_{\Theta}^{\dagger} = h_2 k_{\Theta,k} + k_{\Theta,c}$$



Free transport solver: $\frac{\partial \mathbf{M}}{\partial t} + \nabla \cdot h_1 \mathbf{F} = \mathbf{0}$

Anisotropic Gaussian Velocity Distribution:

$$f(oldsymbol{v}) = rac{lpha_p}{(2\pi|oldsymbol{P}_p|)^{3/2}} \exp\left[-rac{1}{2}(oldsymbol{v}-oldsymbol{U}_p)\cdotoldsymbol{P}_p^{-1}\cdot(oldsymbol{v}-oldsymbol{U}_p)
ight]$$



Granular Stress Tensor Transport



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$$\frac{\partial \rho_p \alpha_p \sigma_p}{\partial t} + \nabla \cdot \rho_p \alpha_p \left(h_2 \boldsymbol{U}_p \otimes \boldsymbol{\sigma}_p - \frac{2}{3} k_{\Theta}^* \nabla \otimes \boldsymbol{\sigma}_p \right) = \rho_p \alpha_p \left(\boldsymbol{S}_{2,flux} - \boldsymbol{S}_2 \right)$$
$$\boldsymbol{S}_{2,flux} = \frac{2p_p^{\dagger}}{\rho_p \alpha_p} \boldsymbol{S}_p - 2\nu_p^{\dagger} \left[\boldsymbol{S}_p \cdot \nabla \boldsymbol{U}_p + (\nabla \boldsymbol{U}_p)^T \cdot \boldsymbol{S}_p - \frac{2}{3} (\boldsymbol{S}_p : \nabla \boldsymbol{U}_p) \mathbf{I} \right]$$
$$\boldsymbol{S}_2 = \left[\frac{2}{\tau_p} + \frac{(3-e)(1+e)}{2\tau_e} \right] \boldsymbol{\sigma}_p$$



Wall boundary conditions

Hydrodynamic solver:

$$\nu_p^* \frac{\partial U_{p,t}}{\partial x_w} = -h_{2,w} \phi_s \mathcal{V}_w U_{p,t} - \frac{p_{p,f} \tan \phi_w}{\rho_p \alpha_p} \frac{U_{p,t}}{|U_{p,t}|} \qquad \mathcal{V}_w = (\pi/6) \sqrt{3\Theta_p}$$
$$k_{\Theta}^{\dagger} \frac{\partial \Theta_p}{\partial x_w} = h_{2,w} \left[\phi_s \mathcal{V}_w |U_{p,t}|^2 - \frac{3}{2} \left(1 - e_w^2\right) \mathcal{V}_w \Theta_p \right]$$

Free transport solver:

$$egin{aligned} \mathbf{F}_w &= h_{1,w}(lpha_p) \int_{oldsymbol{v}\cdotoldsymbol{n}_w>0} \mathbf{G}_r(oldsymbol{v}) \left(oldsymbol{v}\cdotoldsymbol{n}_w
ight) d\mathcal{S}_w \ f_r(oldsymbol{v}) &= \phi_s f_{r,d}(oldsymbol{v}) + (1-\phi_s) \, f_{r,s}(oldsymbol{v}) \ f_{r,s}(oldsymbol{v}) &= f_i \left(oldsymbol{v}-(1+e_w)(oldsymbol{v}\cdotoldsymbol{n}_w)oldsymbol{n}_w
ight) \end{aligned}$$



Example kinetic theory coefficients in hydrodynamic model for particle phase

$$\begin{split} \tau_{p} &= \frac{4\rho_{p}d_{p}^{2}}{3\rho_{g}\nu_{g}C_{D}Re_{p}} & p_{p,k} = \rho_{p}\alpha_{p}\Theta_{p} \\ Re_{p} &= \frac{\alpha_{g}d_{p}|U_{g}-U_{p}|}{\nu_{g}} & p_{p,c} = 4\rho_{p}\eta\alpha_{p}^{2}g_{0}\Theta_{p} - \rho_{p}\alpha_{p}\nu_{p,b}\nabla\cdot U_{p} \\ C_{D} &= \max\left[\frac{24}{Re_{p}}\left(1+Re_{p}^{0.687}\right), 0.44\right]\alpha_{g}^{-2.65} & p_{p,f} = Fr\left(\frac{\alpha_{p}-\alpha_{p,fr:min}}{(\alpha_{p,max}-\alpha_{p})^{s}}\right)^{r} \\ \eta &= \frac{1}{2}(1+e) & \nu_{p,k} = \frac{1}{2}\Theta_{p}\left[\frac{1}{\tau_{p}} + \frac{\eta(2-\eta)}{\tau_{c}}\right]^{-1}\left[1 + \frac{8}{5}\eta(3\eta-2)\alpha_{p}g_{0}\right] \\ g_{0} &= \frac{1-\frac{1}{2}\alpha_{p}}{(1-\alpha_{p})^{3}} & \nu_{p,c} = \frac{8\eta\alpha_{p}g_{0}}{5}\nu_{p,k} + \frac{3}{5}\nu_{p,b} \\ \tau_{c} &= \frac{d_{p}}{6\alpha_{p}g_{0}\sqrt{\Theta_{p}/\pi}} & \nu_{p,f} = \frac{p_{p,f}}{\rho_{p}\alpha_{p}||S_{p}||}\sin\phi \\ \Delta^{*} &= \eta^{2}\Theta_{p}\mathbf{I} + (1-\eta)^{2}\mathbf{P}_{p} & k_{\Theta,k} = \frac{5}{2}\Theta_{p}\left[\frac{3}{\tau_{p}} + \frac{4\eta(41-33\eta)}{\tau_{c}}\right]^{-1}\left[1 + \frac{12}{5}\eta^{2}(4\eta-3)\alpha_{p}g_{0}\right] \\ k_{\Theta,c} &= \frac{12\eta\alpha_{p}g_{0}}{5}k_{\Theta,k} + \frac{3}{2}\nu_{p,b} \end{split}$$





Solution algorithm







Example results, test case 2 : vertical channel riser

 $< \alpha_p >= 0.01$





Conclusions

- 1. A solution algorithm is proposed to accurately treat all fluid-particle regimes occurring simultaneously.
- 2. This algorithm is based on splitting the free-transport flux solver dynamically and locally in the flow. In close-packed to moderately dense regions, a hydrodynamic solver is employed, while in dilute to very dilute regions a kinetic-based finite-volume solver is used in conjunction with quadrature-based moment methods.
- 3. To illustrate the accuracy and robustness of the proposed solution algorithm, it is implemented for particle velocity moments up to second order, and applied to simulate gravity-driven, gas-particle flows exhibiting cluster-induced turbulence.
- 4. By varying the average particle volume fraction in the flow domain, it is demonstrated that the flow solver can handle seamlessly all flow regimes present in fluid-particle flows.





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Questions ?





Semi-discretized equations















