

Quantifying the Uncertainty of Kinetic Theory Predictions of Clustering

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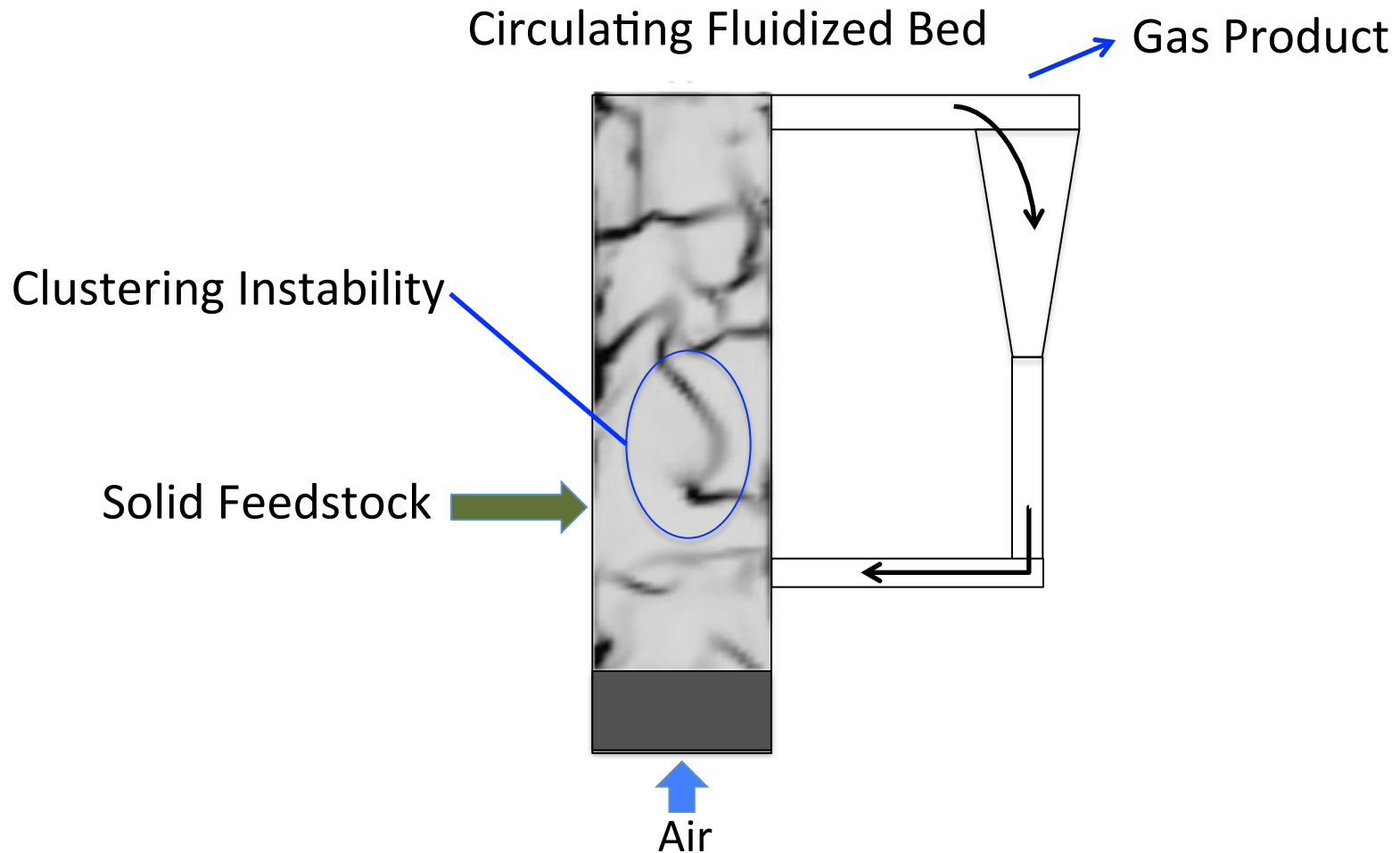


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May 20th, 2013
Pittsburgh, PA
CCR Conference

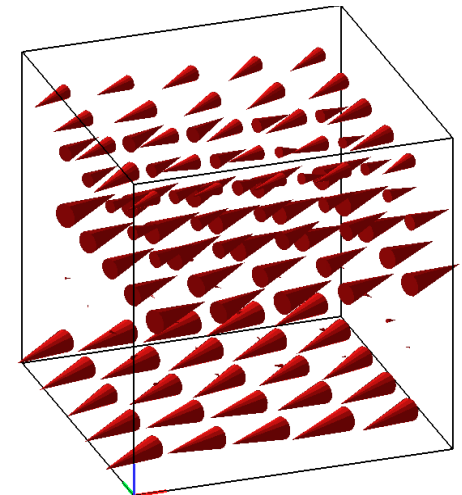
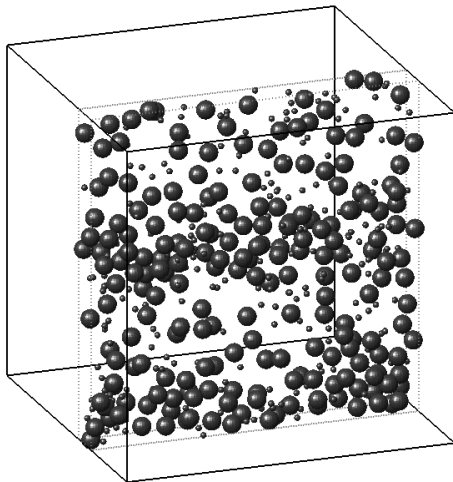
Motivation: Instabilities in Particulate Flows



CFD simulation from Agrawal *et al.* *J. Fluid Mech.* (2001)

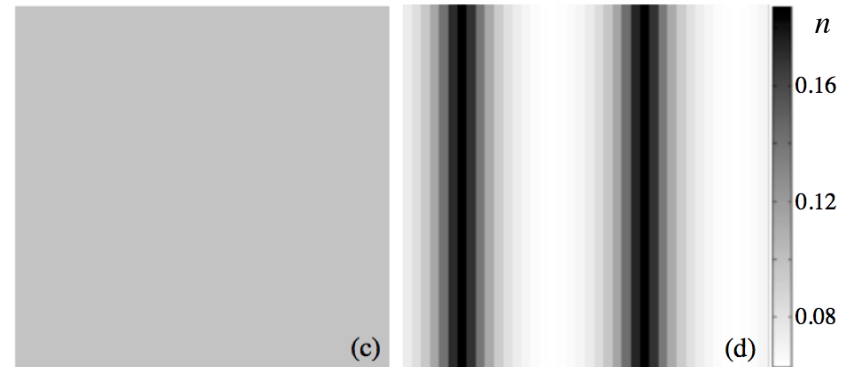
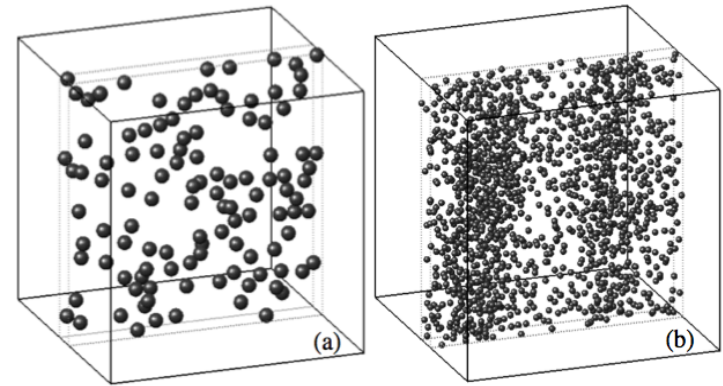
Overall Goal

- (i) Assess kinetic-theory-based continuum models via ability to quantitatively predict instabilities
- (ii) Understand relative importance of clustering mechanisms



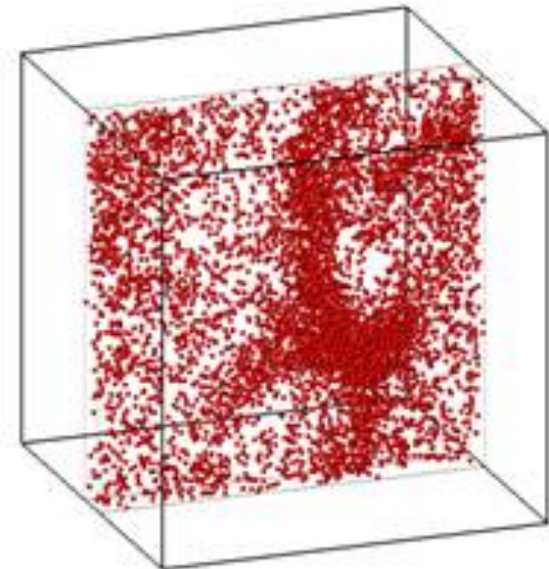
Approach

- Discrete particle simulations (free of kinetic theory)
 - Molecular dynamics (**MD**) →
 - Direct numerical simulation (**DNS**)
- Continuum model (based on kinetic theory)
 - Linear stability analysis
 - Transient continuum simulations →



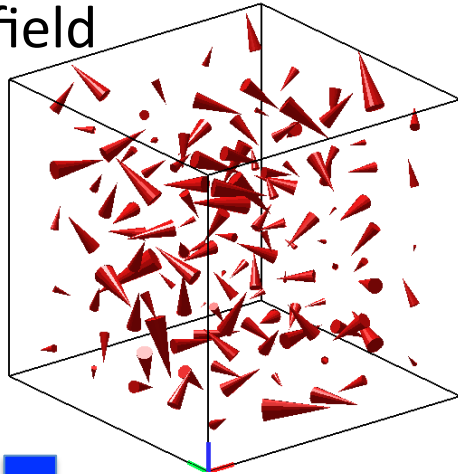
Homogeneous cooling system (HCS)

- System properties
 - No external forces
 - Periodic boundaries in all 3 directions
 - No gradients in the hydrodynamic variables
- Particle properties
 - Constant coefficient of restitution (e)
 - Spherical particles of concentration ϕ
 - Monodisperse, frictionless



Background

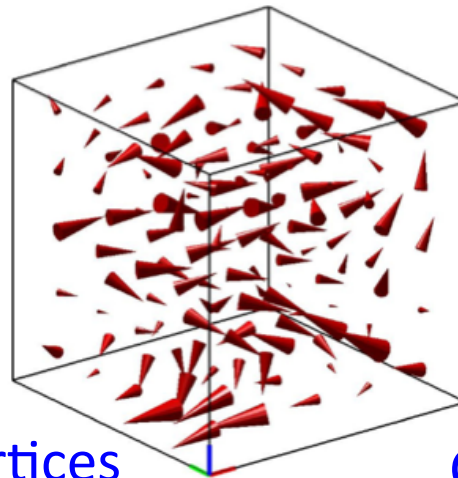
Velocity field



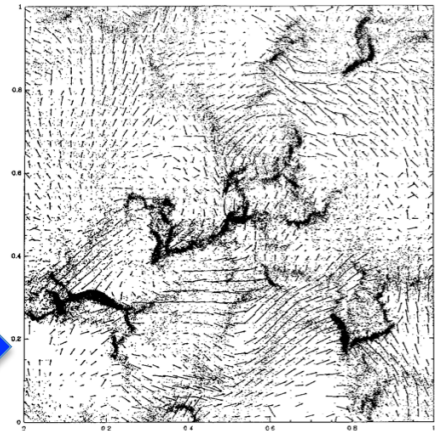
Molecular dynamics (MD) simulations of the HCS

- Dissipative collisions
- Sufficiently large system domain

Velocity field



Particle locations



Vortices

Clusters

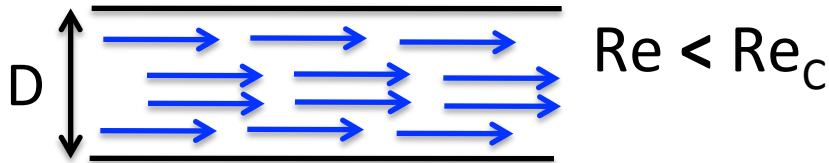
Goldhirsch *et al.*, *J. Sci. Comput.* (1993)

Onset of Instability: A Dimensionless Length Scale

Fluid Flow

$$\text{Re} = \frac{\rho v D}{\mu}$$

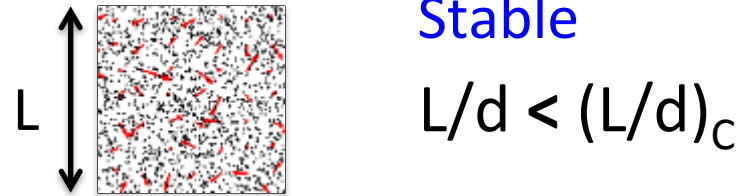
Stable (laminar)



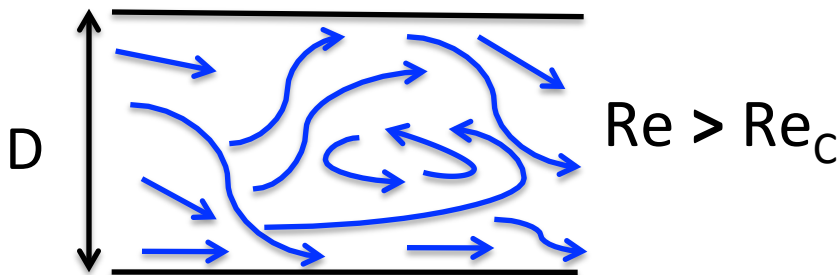
Granular Flow

$$\frac{L}{d} \equiv \frac{\text{Linear domain size}}{\text{particle diameter}}$$

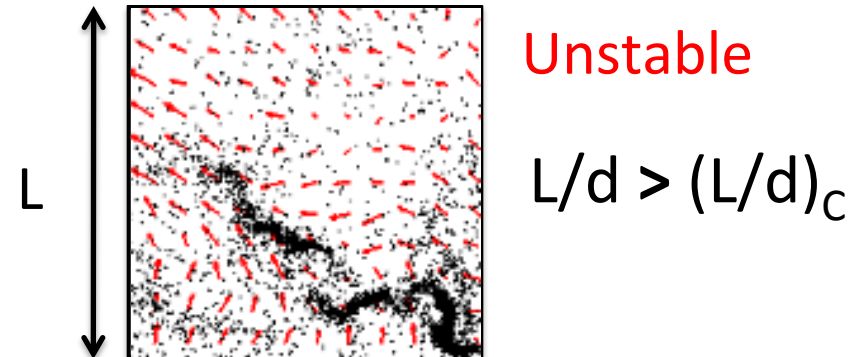
Stable



Unstable (turbulent)

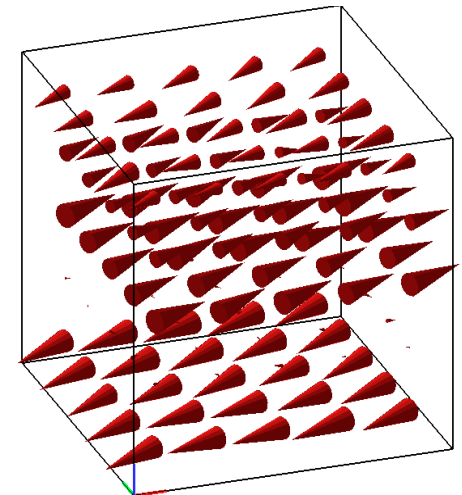
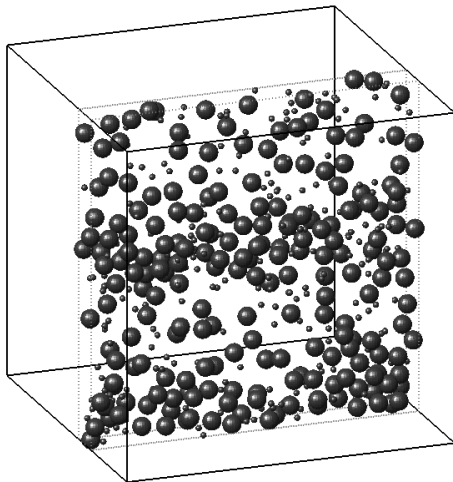


Unstable



Overall Goal

- (i) Assess kinetic-theory-based continuum models via ability to quantitatively predict instabilities



The continuum model

Mass balance $\frac{\partial n}{\partial t} + \underline{u} \cdot \nabla n + n \nabla \cdot \underline{u} = 0$

Momentum balance $\rho \left[\frac{\partial \underline{u}}{\partial t} + \underline{u} \cdot \nabla \underline{u} \right] = -\nabla \underline{\underline{P}}$

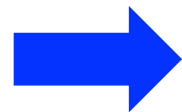
Energy balance $\frac{\partial T}{\partial t} + \underline{u} \cdot \nabla T = -\frac{2}{3n} \left(\nabla \cdot \underline{q} + \underline{\underline{P}} : \nabla \underline{u} \right) - \zeta T$

Chapman-Enskog expansion for heat flux

$$\underline{q} = -\kappa \nabla T - \mu \nabla n - \kappa_2 \nabla T^2 - \mu_2 \nabla n^2 + \dots$$

Truncate at first order in gradients (Navier-Stokes-order)

Thermal conductivity



New coefficient for solids flow (small Knudsen number)

(Granular) Pressure tensor

Heat flux

Cooling rate

Second-order coef.

The continuum model

Mass balance $\frac{\partial n}{\partial t} + \underline{u} \cdot \nabla n + n \nabla \cdot \underline{u} = 0$

Momentum balance $\rho \left[\frac{\partial \underline{u}}{\partial t} + \underline{u} \cdot \nabla \underline{u} \right] = -\nabla \underline{\underline{P}}$

Energy balance $\frac{\partial T}{\partial t} + \underline{u} \cdot \nabla T = -\frac{2}{3n} \left(\nabla \cdot \underline{\underline{q}} + \underline{\underline{P}} : \nabla \underline{u} \right) - \zeta T$

(Granular)
Pressure tensor

Heat flux

Cooling rate

Truncate at first order
in gradients

(Navier-Stokes-order)

$$\underline{\underline{q}} = -\kappa \nabla T - \mu \nabla n$$

$$\zeta = \zeta_0 + \zeta \nabla \cdot \underline{u}$$

$$P_{ij} = p \delta_{ij} - \eta \left(\nabla_j u_i + \nabla_i u_j - \frac{2}{3} \delta_{ij} \nabla \cdot \underline{u} \right) - \gamma \delta_{ij} \nabla \cdot \underline{u}$$

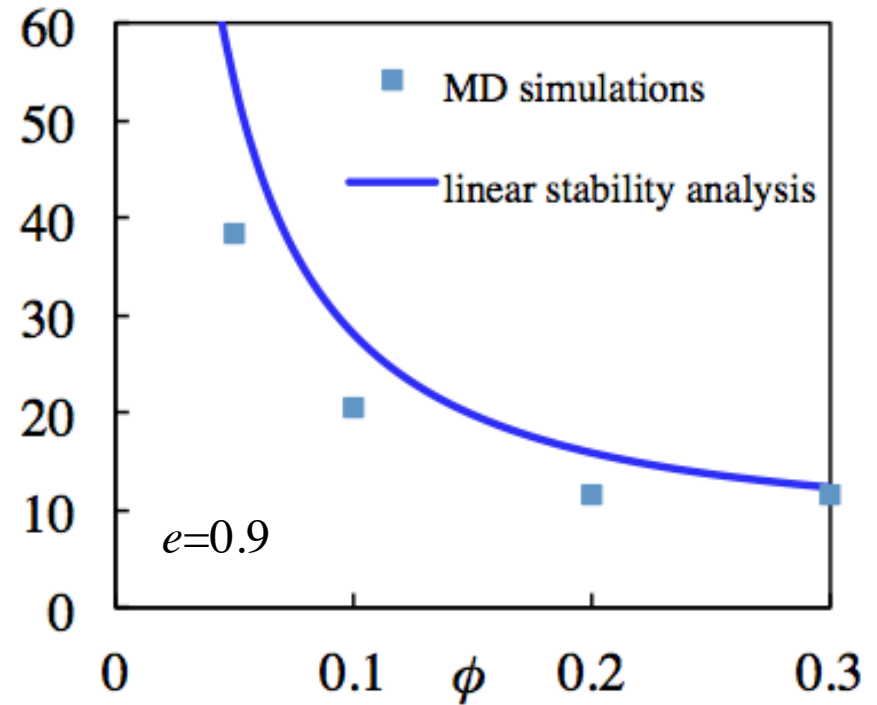
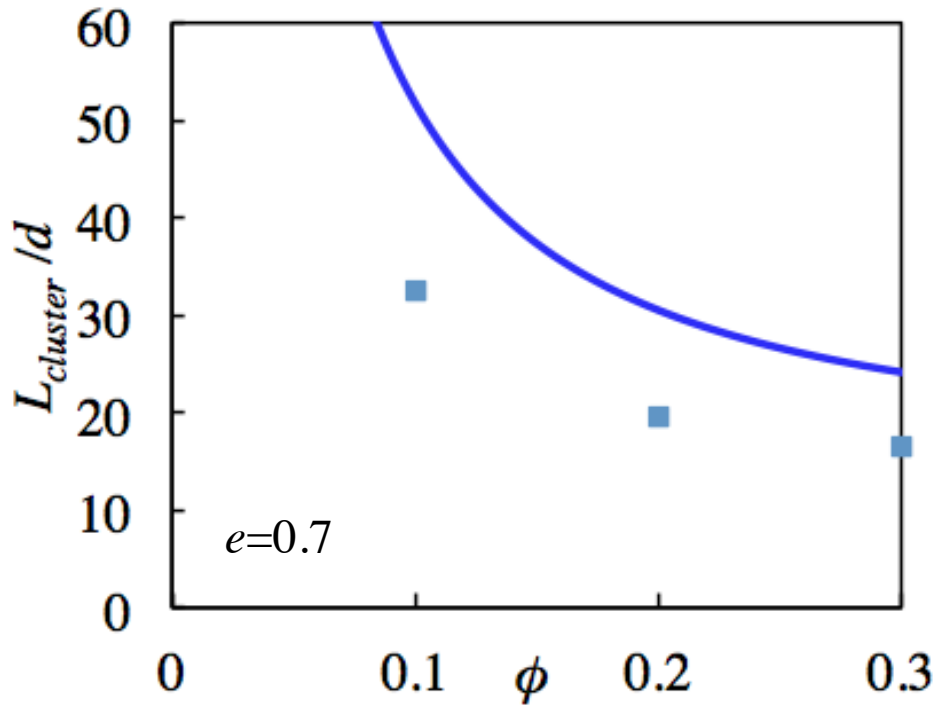
Implies small gradients
(small Knudsen number)

Hydrostatic pressure

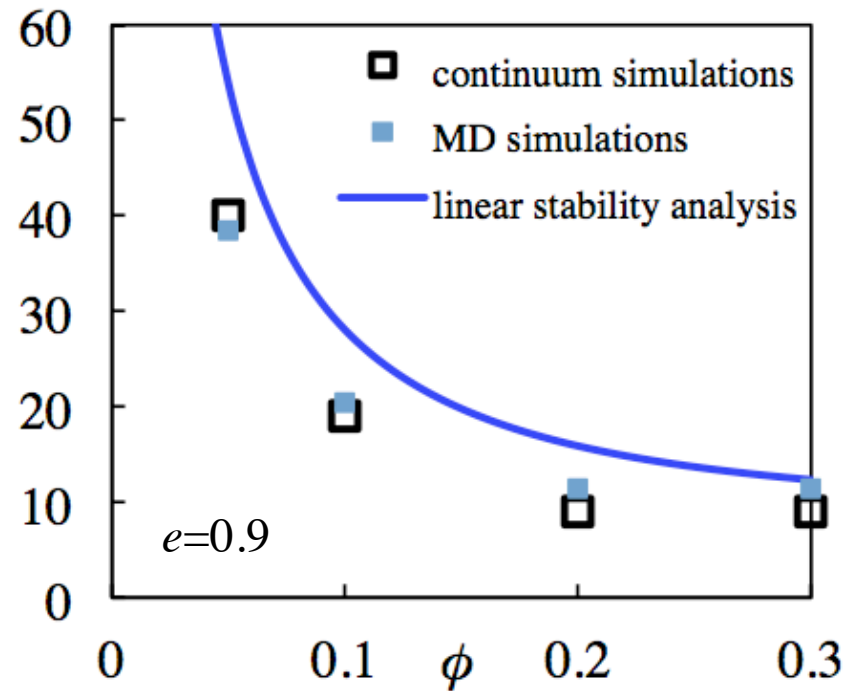
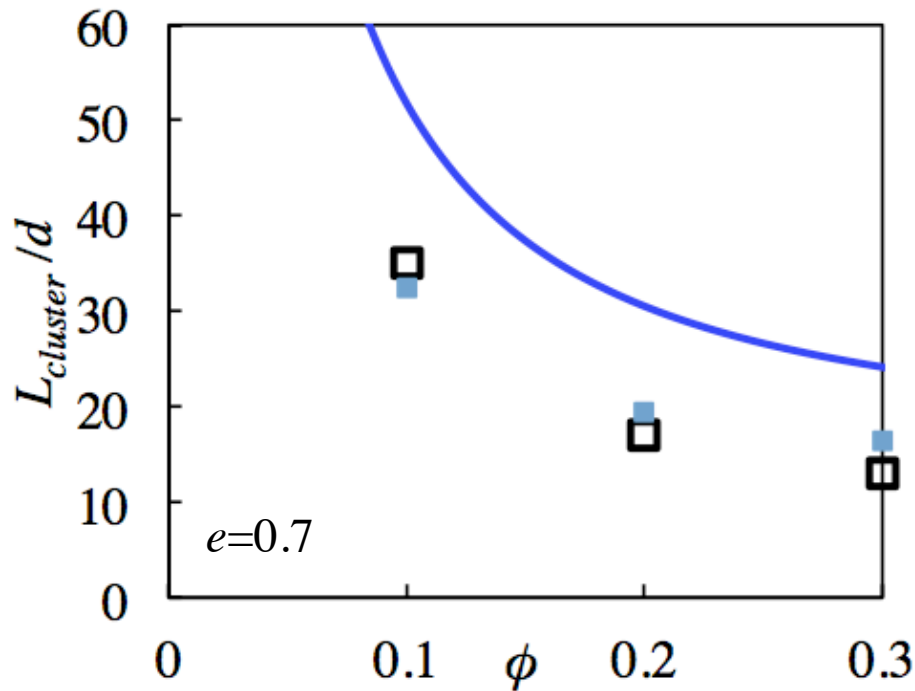
Shear viscosity

Bulk viscosity

High-gradient flows



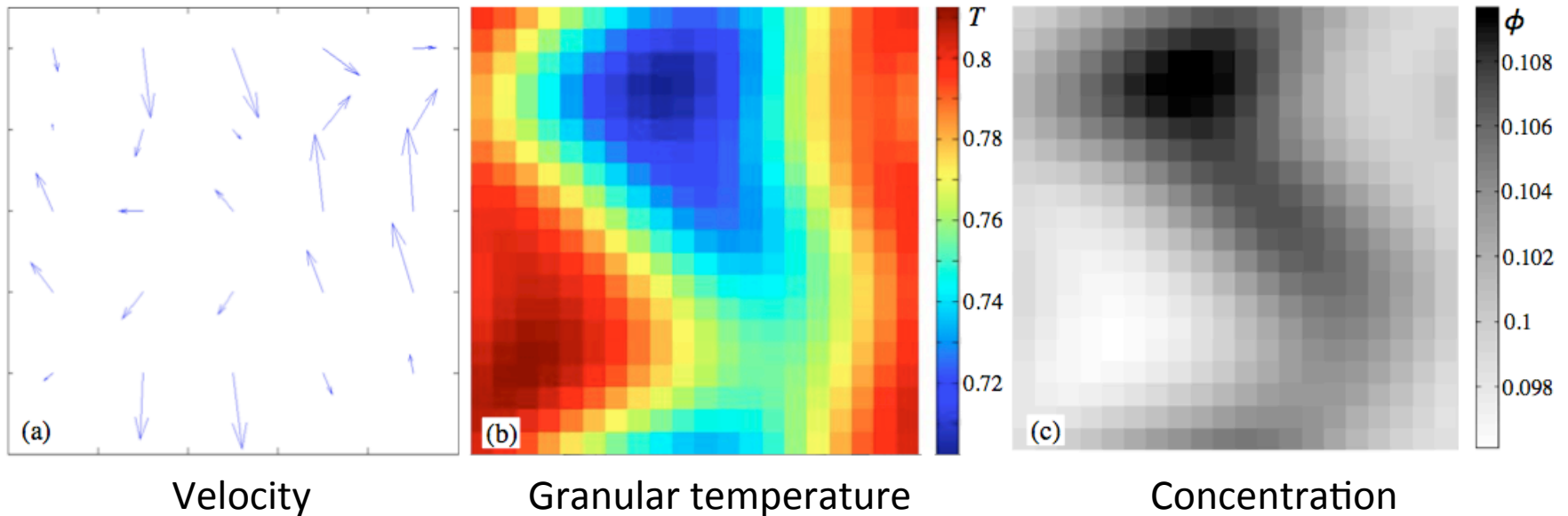
High-gradient flows



Excellent agreement when nonlinearity is considered

High-gradient flows

Coarse-grained hydrodynamic fields at the time of cluster onset



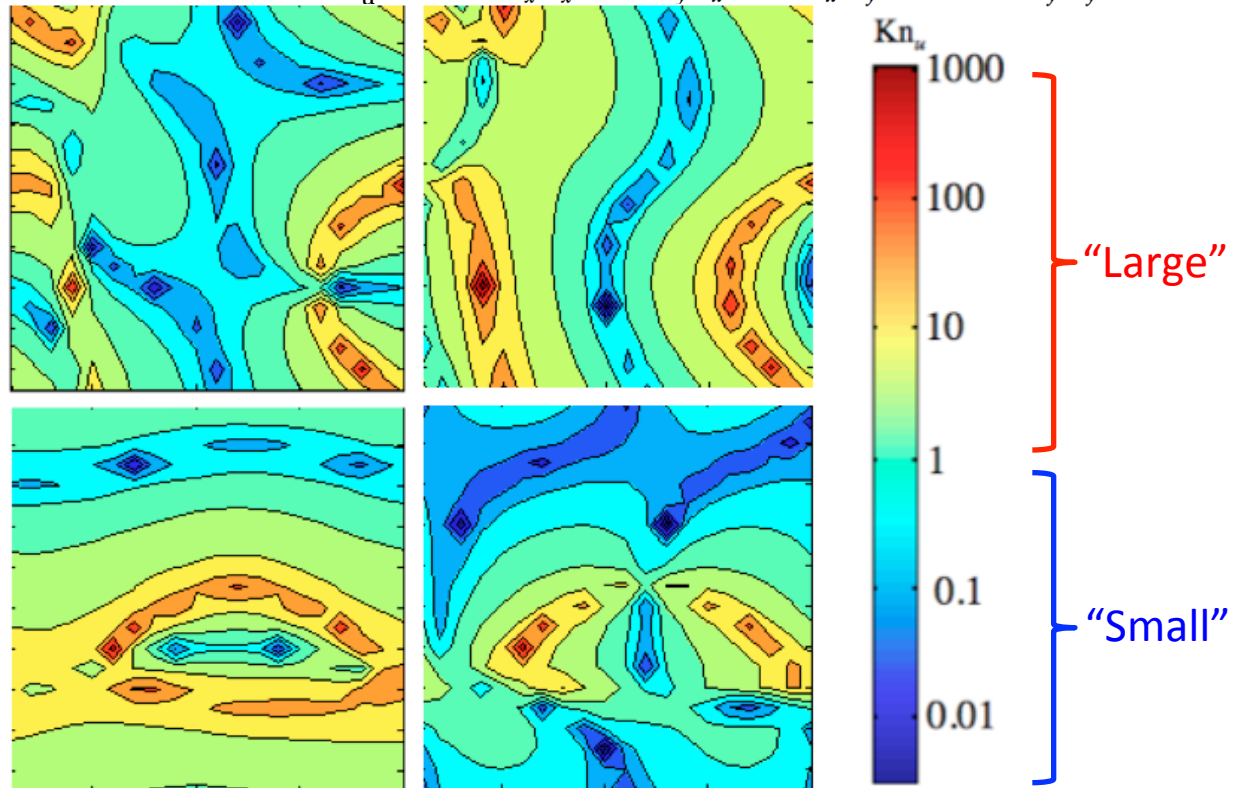
Qualitative evidence of velocity gradients

High-gradient flows

Measure of gradients in **velocity field** at the time of cluster onset

Velocity Knudsen number Kn_u for (a) $\partial_x U_x$, (b) $\partial_y U_x$, (c) $\partial_x U_y$, and (d) $\partial_y U_y$

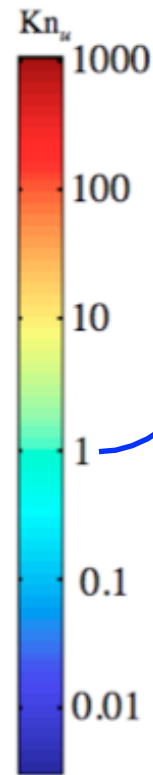
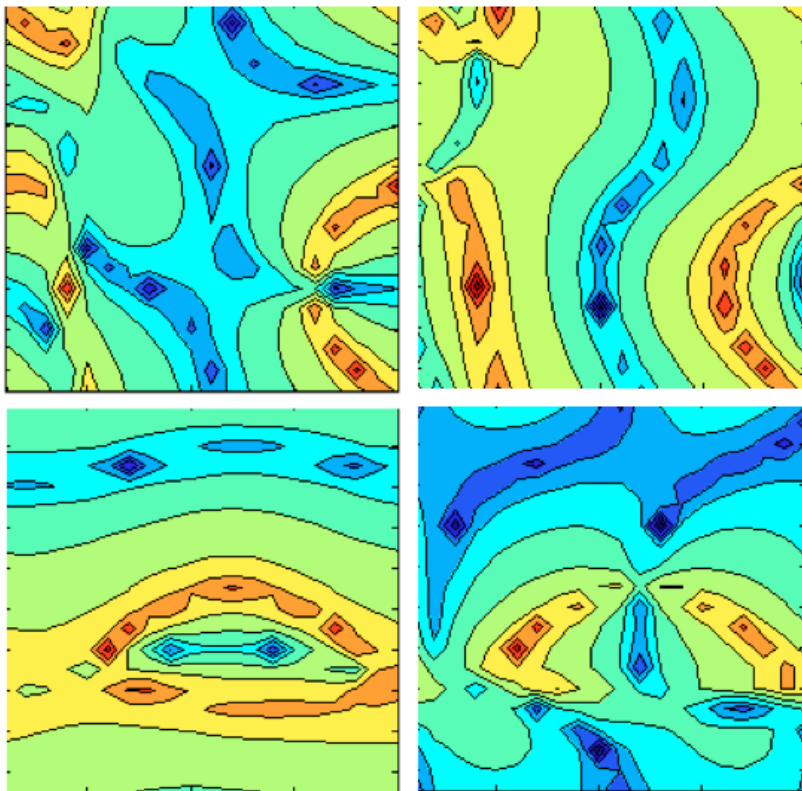
$$Kn_u = \frac{MFP}{L_{gradient}}$$



High-gradient flows

Measure of gradients in **velocity field** at the time of cluster onset

Velocity Knudsen number Kn_u for (a) $\partial_x U_x$, (b) $\partial_y U_x$, (c) $\partial_x U_y$, and (d) $\partial_y U_y$



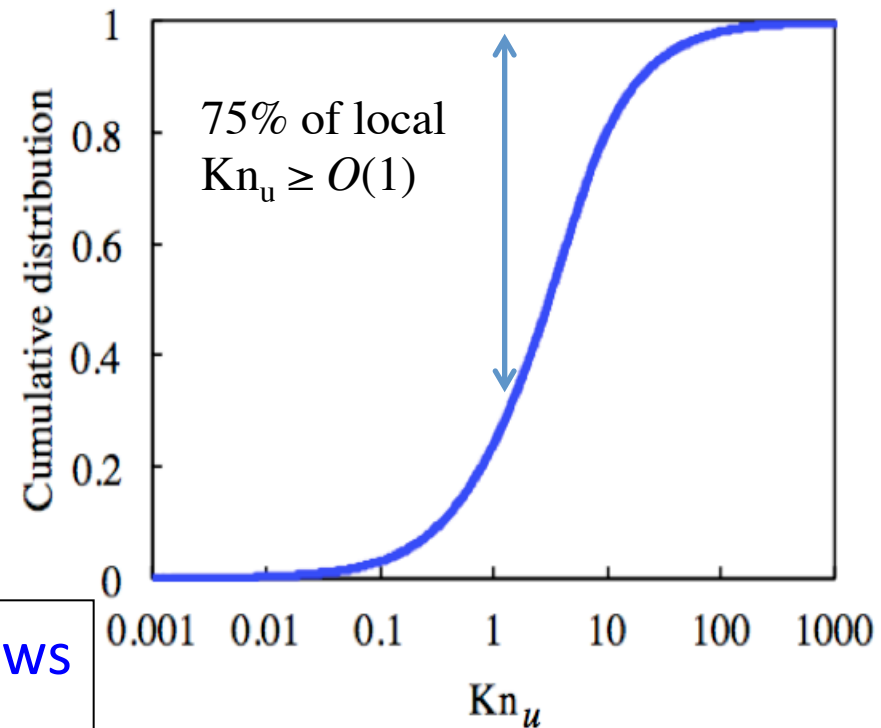
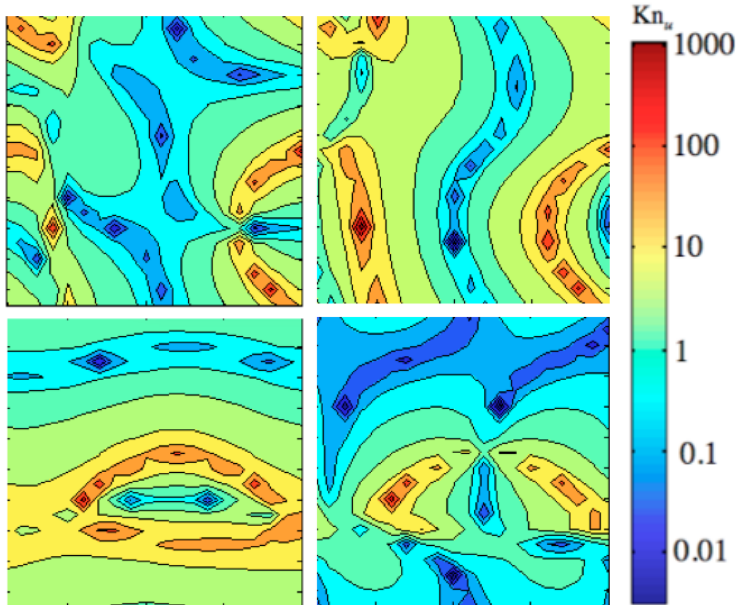
First-order truncation implies $\text{Kn} < \text{Kn}^2$ however $\text{Kn}_u > O(1)$

$$\underline{q} = -\kappa \nabla T - \mu \nabla n - \kappa_2 \nabla T^2 - \mu_2 \nabla n^2 + \dots$$

High-gradient flows

Measure of gradients in **velocity field** at the time of cluster onset

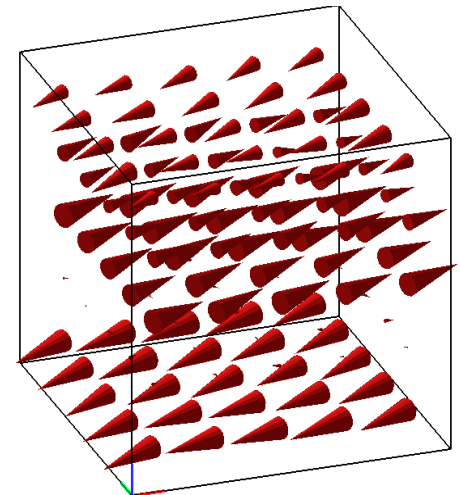
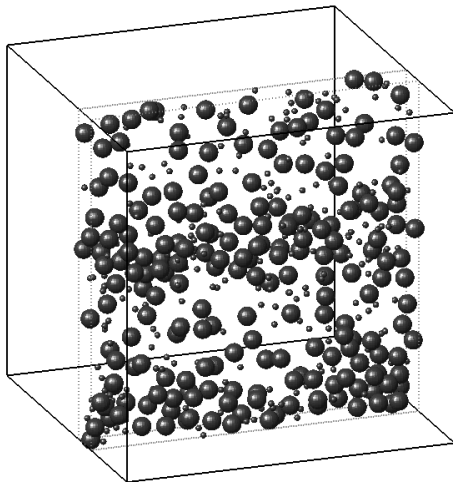
Velocity Knudsen number Kn_u for (a) $\partial_x U_x$, (b) $\partial_y U_x$, (c) $\partial_x U_y$, and (d) $\partial_y U_y$



Theory does well in high-gradient flows despite small-Kn assumption

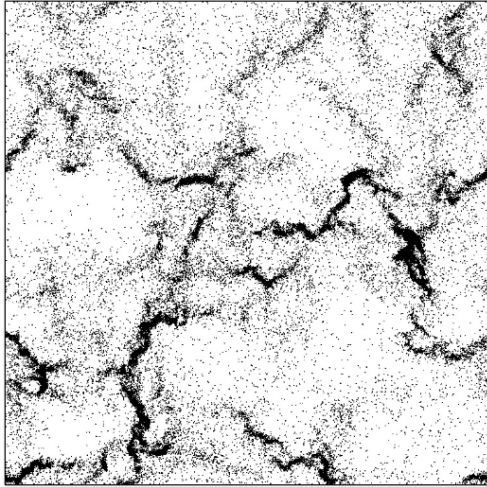
Overall Goal

- (i) Understand relative importance of clustering mechanisms



Particle Clustering Instability: Known Mechanisms

Homogeneous Cooling System



Goldhirsch, *et al.*, *J. Sci. Comput.* (1993)

Granular Work

Solid Effects
(Dissipative Collisions)

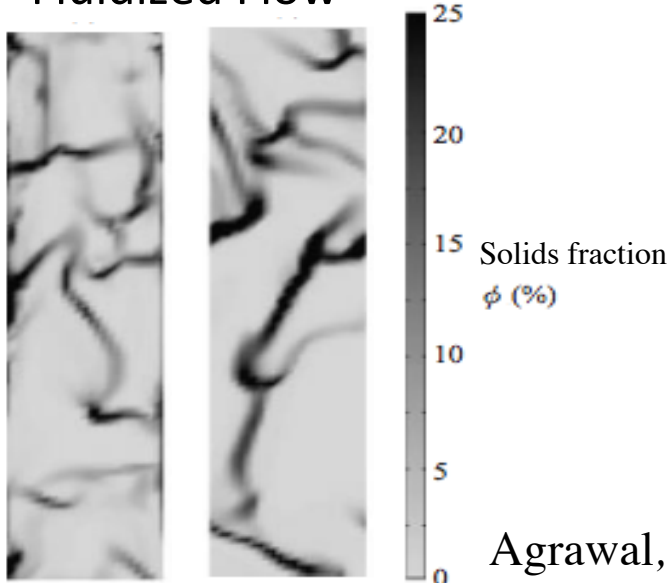
Inelasticity

- Hopkins & Louge 1991
- Goldhirsch *et al.* 1993

Friction

- Mitrano *et al.* 2013

Fluidized Flow



Fluidization Work

Fluid Effects

Mean Drag

- Glasser *et al.* 1998
- Agrawal *et al.* 2001

“Viscous Losses”

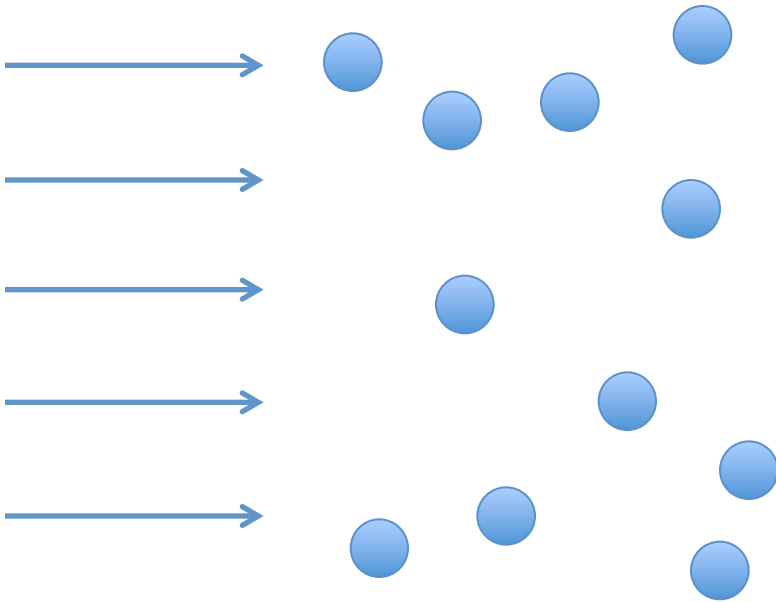
- Wylie & Koch 2000

Agrawal, *et al.*, *J. Fluid Mech.* (2001)

Mean Drag v. Viscous Losses

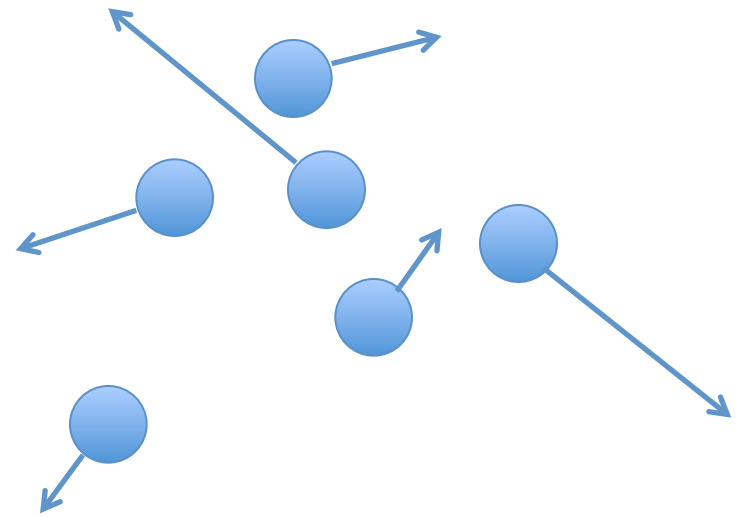
Gas flow

fixed particles



Mean relative velocity decreases
 When $\bar{U}_g - \bar{U}_s = 0$, $F_{fluid} = 0$

Particle flow in gas with $\bar{U}_g - \bar{U}_s = 0$



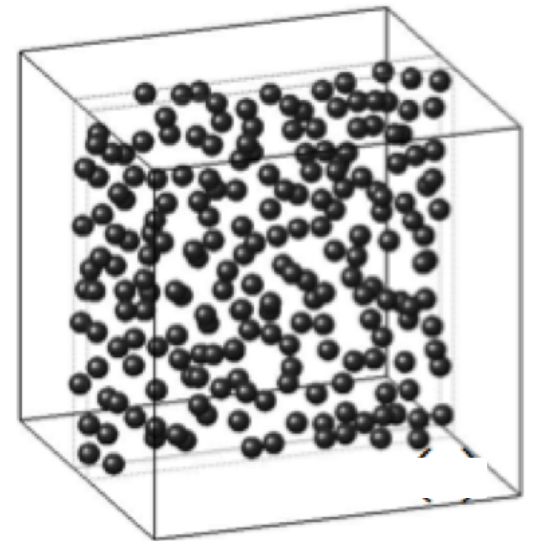
Granular temperature decreases
 When $T = 0$, $F_{fluid} = 0$

Conditions

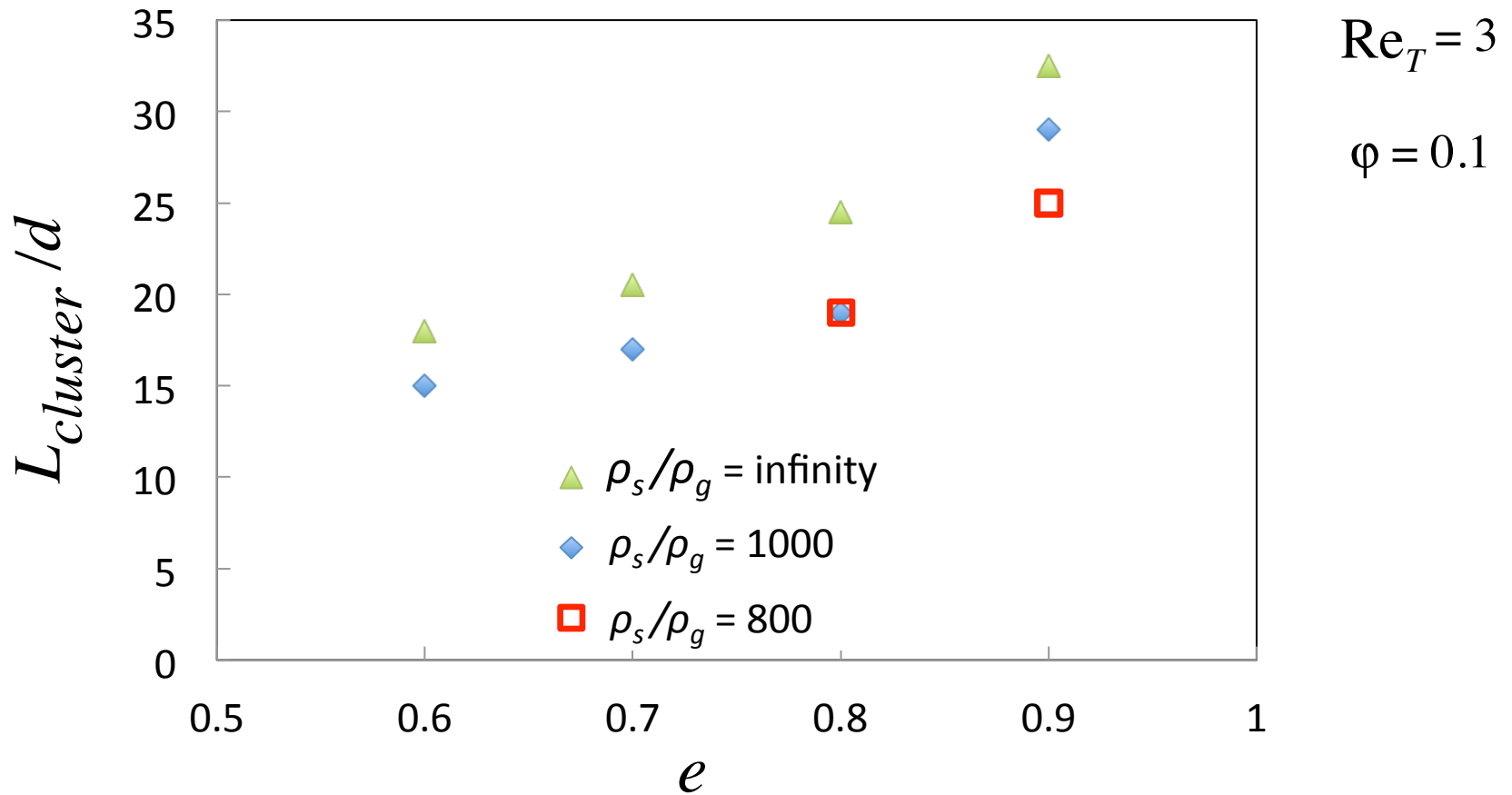
- Restitution coefficient: $0.8 \leq e \leq 1.0$
- Solids fraction: $0.1 \leq \phi \leq 0.4$
- Density ratio: $800 \leq \frac{\rho_{solid}}{\rho_{fluid}} \leq 1500$

$$\text{Re}_T \propto \sqrt{T_0}$$

$$\begin{aligned} \text{Re}_M &\propto V_{rel} \\ &\propto \frac{\text{Particle inertial forces}}{\text{Fluid viscous forces}} \end{aligned}$$

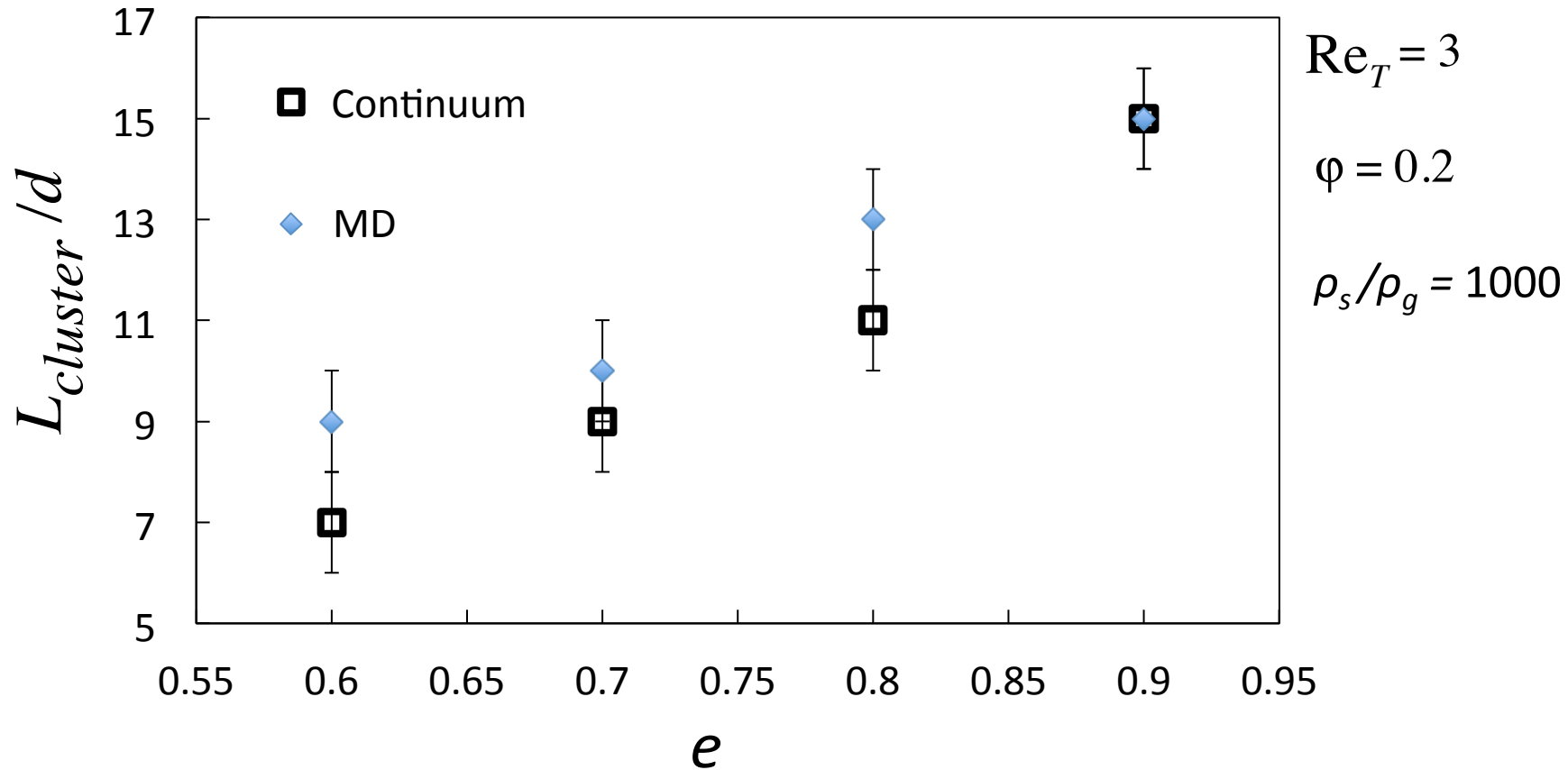


Gas-solid Continuum Model of HCS



Both viscous losses and collisional dissipation important

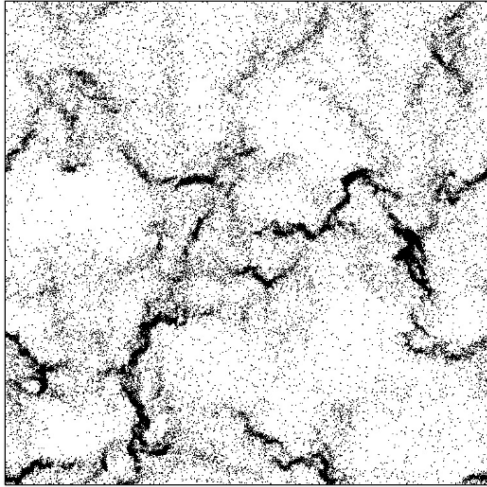
Gas-solid HCS: Continuum v. MD



Strong Agreement between continuum model and MD

Particle Clustering Instability: Known Mechanisms

Homogeneous Cooling System



Goldhirsch, *et al.*, *J. Sci. Comput.* (1993)

Granular Work

Solid Effects
(Dissipative Collisions)

Inelasticity

- Hopkins & Louge 1991
- Goldhirsch et al. 1993

Friction

- Mitrano et al. 2013

Fluidized Flow



Solids fraction
 ϕ (%)

Fluidization Work

Fluid Effects

Mean Drag

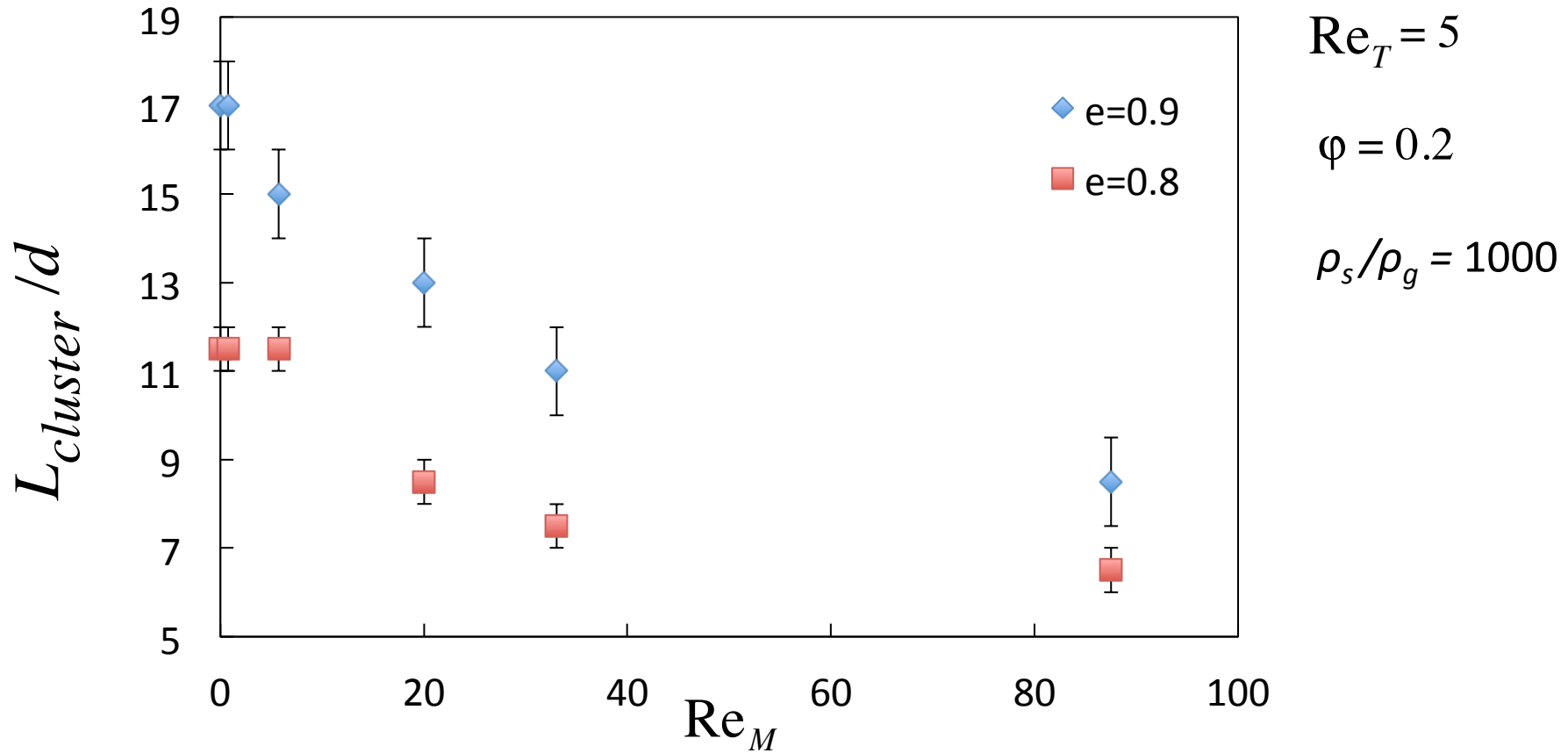
- Glasser et al. 1998
- Agrawal et al. 2001

Viscous Losses

- Wylie & Koch 2000

Agrawal, *et al.*, *J. Fluid Mech.* (2001)

Gas-solid Continuum Model of Settling Flow



Both mean drag and collisional dissipation important

Concluding Remarks

- Small Knudsen number assumption not so restrictive
- Collisional dissipation, mean drag, and viscous losses all important in conditions studied
- Strong agreement between continuum model and discrete particle simulations

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